

Characterization of vacancy-type defects in Mg- and N-implanted GaN by using a monoenergetic positron beam

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Abstract—Annealing behaviors of vacancy-type defects in Mg and N-implanted GaN were studied by positron annihilation. The major defect species in as-implanted samples was identified as Ga-vacancy (V_{Ga})-type defects. For Mg-implanted GaN with sequential N-implantation after annealing above 1000°C, the defect species were vacancy clusters such as $(V_{\text{Ga}}V_{\text{N}})_3$. Due to the downward shift of the Fermi level position resulting from a partial activation of Mg, the charge states of defects tended to become positive. For N-implanted GaN, the size of the vacancy cluster started to decrease above 1200°C annealing, which was attributed to recombinations between V_{NS} coupled with V_{Ga} s and excess N atoms. The impact of sequential N-implantations on vacancies in Mg-implanted GaN was found to be most pronounced when the ratio of the concentration of N to that of Mg was three.

Index Terms—GaN, ion implantation, vacancy defects, positrons

I. INTRODUCTION

Power devices are solid-state electronics that control and convert electrical energy. GaN is a key material for both present and future power electronics because of its excellent physical properties [1]. Currently, nearly all commercial GaN-based devices are fabricated using a lateral

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architecture. Due to the power-handling limitations of the lateral devices, however, vertical GaN power devices have garnered significant attention [2–4]. The fabrication of vertical devices requires a precisely controlled impurity doping technique, and this can be achieved using ion implantation. One of the obstacles to fabricating vertical devices has been the activation of an implanted p-type impurity: Mg. Recently, it was reported that ultra-high-pressure annealing (UHPA) is effective in obtaining a high activation rate of Mg implanted into GaN [4–6]. One of the reasons for the difficulties in activating implanted Mg is suggested to be the introduction of donor-like defects such as nitrogen vacancies (V_{NS}) [7]. Thus, the sequential implantation of N after Mg-implantation is also an effective technique to increase the Mg activation rate [4,8]. However, the reactions of defects and the formation of secondary defects during UHPA, as well as the impact of sequential N-implantation on them, are not well known. Positron annihilation spectroscopy is a useful technique for detecting vacancy-type defects in solids [9,10]. This technique has been used to study native and process-induced defects in GaN [11–15]. In the present study, we used a monoenergetic positron beam to study the annealing behaviors of vacancy-type defects in Mg- and N-implanted GaN. Photoluminescence (PL) measurements were carried out in order to correlate vacancy-type defects with non-radiative recombination centers (NRC). Secondary defects were characterized by low-angle annular dark-field scanning transmission electron microscopy (LAADF-STEM).

II. EXPERIMENTAL

The present samples were undoped GaN grown on GaN substrates by metal-organic vapor phase deposition. Mg ions were implanted with energies ranging from 10 to 700 keV in order to obtain a 700-nm-deep box profile. The concentration of Mg ($[\text{Mg}]$) in the box profile was $1 \times 10^{18} \text{ cm}^{-3}$. After the Mg-implantation, N ions were implanted, and their depth profile was set to be similar to that of Mg. The concentrations of N in the box profiles were varied at 1×10^{18} , 3×10^{18} , and $1 \times 10^{19} \text{ cm}^{-3}$. After the implantation, the samples were annealed at temperatures up to 1450°C with an annealing time (t_a) of 5 minutes under a N_2 pressure (P_{N_2}) of 1 GPa using a UHPA system [6]. UHPA with $P_{\text{N}_2}=0.5$ GPa was also performed with $t_a=5\text{--}60$ min.

LAADF-STEM was used to examine the microstructure

using a Titan G2 80-200 (Thermo Fisher Scientific) microscope at an accelerating electron voltage of 200 kV. The samples were prepared using a focused-ion-beam (FIB) scanning electron microscope, the dual-beam Helios 5UX (Thermo Fisher Scientific). The sample thickness was 30 nm based on the intensity of backscattered electrons [16,17]. The depth distributions of Mg were measured by secondary ion mass spectrometry. PL spectra were measured at 13 K using a 325-nm He-Cd laser as an excitation source under a weak-excitation condition (38 Wcm^{-2}).

Vacancy-type defects in the samples were characterized using a monoenergetic positron beam. Details on this technique can be found elsewhere [9,10,12–14]. Doppler broadening spectra were measured using Ge detectors as a function of the incident positron energy (E), and they were characterized by the S parameter. The spectra were also measured using a coincidence Doppler broadening technique [9]. The relationship between S and E was analyzed using the computer code VEPFIT [18]. The spectra were measured in darkness and under the illumination of a 325-nm He-Cd laser (Kinmon Koha, IR3802R-G). The Doppler broadening spectra were simulated by QMAS (Quantum MATERIALS Simulator) code [19] using the projector augmented-wave (PAW) method [20] and a plane-wave basis set. Details on applying QMAS to calculations of positron states and annihilation parameters in group-III nitrides are given elsewhere [21].

III. RESULTS AND DISCUSSION

Figure 1 shows the S parameter of Mg-implanted GaN with sequential N-implantation ($[N]=1 \times 10^{18} \text{ cm}^{-3}$) before and after UHPA ($P_{N_2}=1 \text{ GPa}$) as a function of the incident positron energy E [15]. The mean implantation depth of positrons is shown on the upper horizontal axis. The S value corresponding to the annihilation of positrons in defect-free GaN (S_f) was 0.442 [12–14]. For the as-implanted sample, the observed S values were higher than S_f , indicating the trapping of positrons by vacancy-type defects. The S values for the damaged region ($E=1\text{--}15 \text{ keV}$) were increased by annealing after 1000°C , but they started to decrease above 1100°C annealing. For the samples annealed above 1300°C , the S value was close to S_f , suggesting that the trapping of positrons by vacancies is almost negligible. The solid curves shown in Fig. 1 are fits to the experimental data. Figure 2 shows the derived depth distributions of S and Mg. No significant change in the Mg profiles was observed up to 1300°C annealing. For the sample annealed at 1000°C , the increase in S occurred below 700 nm, and this region agreed with the box profile of Mg, suggesting that the agglomeration of vacancy-type defects mainly occurred in the region with high [Mg].

The annealing behaviors of the S values for Mg-implanted GaN with $[Mg]$ of $1 \times 10^{17}\text{--}1 \times 10^{19} \text{ cm}^{-3}$ were reported by ref. 22. After annealing at 1300°C , the S value increased as $[Mg]$ decreased. This suggests that the primary origin of the decrease in S is a partial activation of Mg by high-temperature annealing, resulting in a decrease of the Fermi level position. The downward shift of the Fermi level position causes a

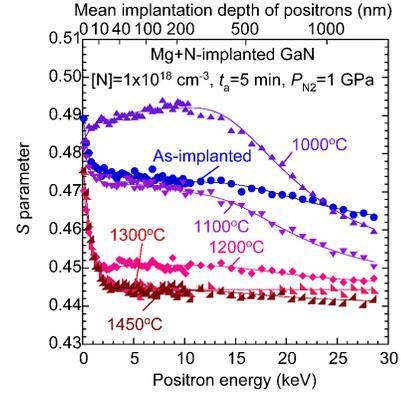


Fig. 1. S parameters as function of incident positron energy E for Mg-implanted GaN with sequential N-implantation before and after annealing at $1000\text{--}1450^\circ\text{C}$. Solid curves are fits to experimental data.

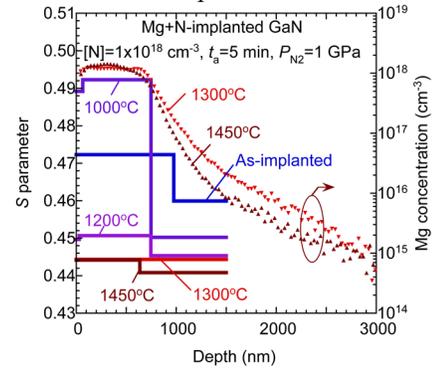


Fig. 2. Depth distributions of S obtained from analysis of S – E curves shown in Fig. 1 (solid lines) and those of Mg after annealing at 1300°C and 1450°C (triangles).

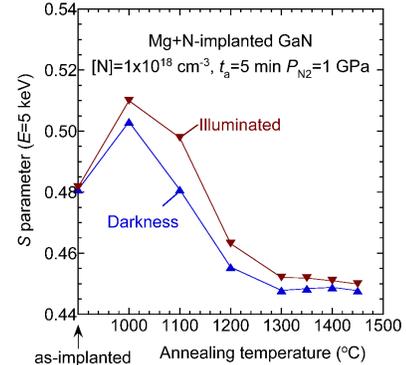


Fig. 3. S parameter measured at $E=5 \text{ keV}$ as function of annealing temperature for Mg-implanted GaN with sequential N-implantation. Measurements were performed in darkness and under He-Cd illumination.

change in the charge state of vacancies from neutral to positive ($V^0 \rightarrow V^+$). The trapping rate of positrons for a positively charged vacancy is smaller than that for a neutral or negatively charged vacancy [9]. Thus, the observed decrease in the S value for the annealed samples (Figs. 1 and 2) can be attributed to the downward shift of the Fermi level position by high-temperature annealing.

Figure 3 shows the annealing behavior of the S value measured by the coincidence Doppler broadening technique at

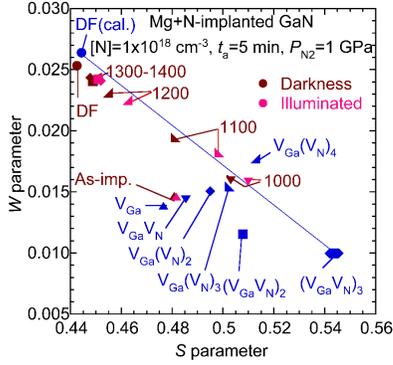


Fig. 4. S – W plots for Mg-implanted GaN with sequential N-implantation before (As-imp.) and after annealing (1000–1400°C). Measurements were done in darkness (brown symbols) and under He-Cd illumination (pink symbols). (S, W) value for defect-free GaN is shown as “DF.” (S, W) values obtained by simulations for defect-free GaN [$D_{F(\text{cal.})}$] and typical vacancy-type defects are shown as blue symbols.

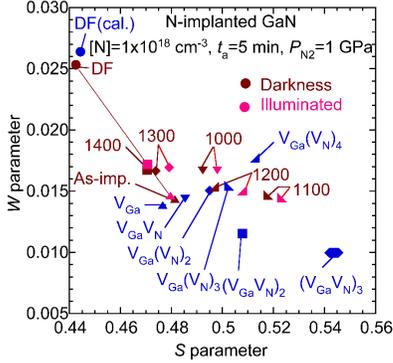


Fig. 5. S – W plots for N-implanted GaN measured in darkness and under He-Cd illumination.

$E=5$ keV ($[N]=1 \times 10^{18}$ cm $^{-3}$ and $P_{N_2}=1$ GPa) [15]. The spectra were also measured under the illumination of a He-Cd laser. The difference between the S values measured in darkness and under illumination can be attributed to the change in the charge state of vacancy-type defects from positive to neutral by the capturing of excited electrons by illumination ($V^+ + e^- \rightarrow V^0$) [12–14]. The illumination effect tended to decrease after annealing above 1200°C. This can be attributed to the suppression of electron capture by vacancies resulting from the activation of Mg by a higher annealing temperature.

Figures 4 and 5 show S – W relationships for the Mg-implanted GaN with sequential N-implantation and N-implanted GaN ($[N]=1 \times 10^{18}$ cm $^{-3}$ and $P_{N_2}=1$ GPa) calculated using the coincidence Doppler broadening spectra, respectively [15]. During the measurements, the value of E was fixed at $E=5$ keV. The (S, W) value corresponding to the positron annihilation in a defect-free GaN is shown as “DF” (defect-free). The (S, W) values measured in darkness and under illumination are shown as brown and pink symbols, respectively. The (S, W) value calculated from the simulated spectrum for the delocalized state is shown as “DF(cal.)” The simulated (S, W) values for typical vacancy-type defects are also shown in the figures (blue symbols).

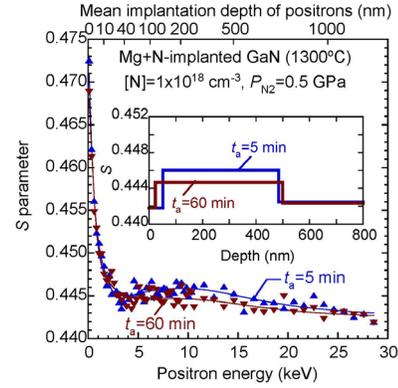


Fig. 6. S – E relationships for Mg-implanted GaN with sequential N-implantation. Annealing temperature was 1300°C, and annealing times (5 and 60 min) are shown in figure. Inset shows depth distributions of S obtained from analysis.

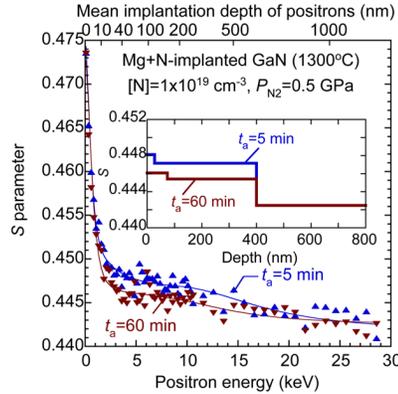


Fig. 7. S – E relationships for Mg-implanted GaN with sequential N-implantation. Inset shows depth distributions of S obtained from analysis.

For the as-implanted samples (As-imp.), the (S, W) values were close to the values for V_{Ga} and $V_{Ga}V_N$. Thus, the major defect species can be identified as V_{Ga} -type defects. For the Mg-implanted GaN with sequential N-implantation (Fig. 4), the (S, W) values for the samples annealed above 1000°C were located along the line connecting the (S, W) values for defect-free GaN and $(V_{Ga}V_N)_3$, suggesting that the major defect species was vacancy clusters. Above 1100°C, the (S, W) value approached the defect-free value, which is attributed to the downward shift of the Fermi level position as discussed above.

For N-implanted GaN after annealing at 1000–1100°C (Fig. 5), the value shifted to the right-hand side, suggesting the agglomeration of vacancy-type defects. After 1200°C annealing, however, the (S, W) value shifted towards the left-hand side. From the observed annealing behavior of the (S, W) value, it can be concluded that the shrinkage of vacancy-type defects starts above 1200°C, and this can be attributed to recombinations between vacancy clusters and excess N atoms. The result also suggests that sequential N-implantation effectively suppresses the agglomeration of vacancies; however, the sample must be annealed above 1200°C to achieve this effect. In Figs. 4 and 5, the (S, W) values tended to show a right-hand shift under the illumination. The increase in

the number of V_N coupled with V_{Ga} caused the right-hand change in the S - W plot. Thus, the vacancy clusters acting as electron trapping centers are considered to have more V_N than neutral or negatively charged clusters detected in darkness.

Figures 6 and 7 show the S - E relationships for Mg-implanted GaN with sequential N-implantation with different $[N]$ s, (1×10^{18} and 1×10^{19} cm^{-3}), respectively ($P_{N_2}=0.5$ GPa, $t_a=5$ and 60 min). The depth distributions of S obtained from the analysis are shown in the insets in Figs. 6 and 7. The S value in the damaged region decreased as the annealing time increased. Figure 8 shows the annealing time dependence of the S value measured for Mg-implanted GaN with sequential N-implantation ($[N]=1 \times 10^{18}$, 3×10^{18} , and 1×10^{19} cm^{-3}). These S values were calculated from coincidence Doppler broadening spectra measured at $E=8$ keV. For both measurements in darkness and under illumination, the S values for the sample with $[N]=3 \times 10^{18}$ cm^{-3} were smaller than those for the samples with $[N]=1 \times 10^{18}$ and 1×10^{19} cm^{-3} , suggesting that the residual vacancy concentration after annealing at 1300°C was low when the ratio of $[N]$ to $[Mg]$ was three.

Figure 9 shows PL spectra for Mg-implanted GaN with sequential N-implantation ($[N]=1 \times 10^{18}$, 3×10^{18} , and 1×10^{19} cm^{-3}). The annealing temperature, annealing time, and N_2 pressure were 1300°C , 60 min, and 0.5 MPa, respectively. The sharp luminescence peak at 3.5 eV corresponds to the recombination of excitons bound to a Mg_{Ga} acceptor (acceptor bound exciton: ABE), and the origin of a broad ultraviolet luminescence (UVL) band (2.6–3.4 eV) is a free electron or a shallow donor to a Mg_{Ga} acceptor transition [23,24]. Thus, the increases in the intensities of the ABE peak and the UVL band relate to the increase in the concentration of Mg_{Ga} and the decrease in the concentration of defects that act as NRCs. As shown in the inset of Fig. 9, the highest ABE peak intensity was obtained for the sample with $[N]=3 \times 10^{18}$ cm^{-3} , which can be associated with the observed suppression in the concentration of residual vacancies in this sample (Fig. 8).

For the sample with $[N]=1 \times 10^{18}$ cm^{-3} , a broad luminescence band at 2.1–2.6 eV was observed, but it was suppressed for the samples with $[N]=3 \times 10^{18}$ and 1×10^{19} cm^{-3} . This green luminescence (GL) band is often observed for Mg-implanted GaN [24] and assigned to a transition involving V_N . Thus, the suppression of the GL band is due to the recombination of V_N and N by sequential N-implantation. However, the total luminescence intensity decreased due to the introduction of NRCs for the sample with $[N]=1 \times 10^{19}$ cm^{-3} .

Figure 10 shows images of the (a) Mg-implanted GaN with sequential N-implantation and (c) N-implanted GaN after annealing at 1400°C obtained by LAADF-STEM [15]. Their magnified images are shown in (b) and (d), respectively. Bright dots were observed in the N-implanted sample. Their origin has been attributed to nano-scale intrinsic defects; an extra Ga atomic plane inserted between the Ga planes of the matrix [25,26]. For the Mg-implanted sample with N-implantation, both bright dots and circles were observed. The nature of the latter defects was identified as collapsed vacancy disks forming intrinsic dislocation loops [25,26]. Since the

formation of such dislocation loops requires the removal of both Ga and N atoms from the inside of the loops, the V_{Ga} -rich condition for the N-implanted GaN can be the origin of the suppression of the formation of such secondary defects.

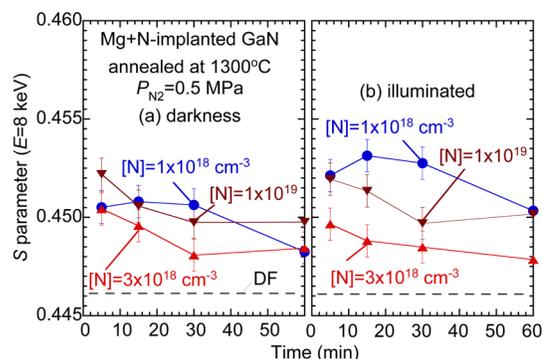


Fig. 8. S parameters as function of annealing time for Mg-implanted GaN with sequential N-implantation ($[N]=1 \times 10^{18}$, 3×10^{18} , and 1×10^{19} cm^{-3}). Annealing temperature was 1300°C . Measurements were performed in (a) darkness and (b) under illumination.

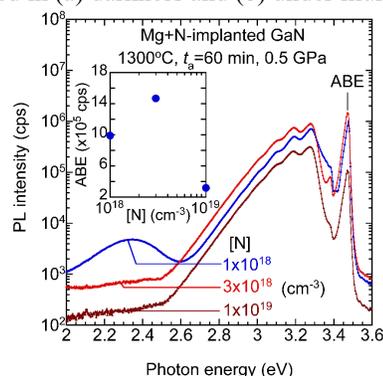


Fig. 9. PL spectra for Mg-implanted GaN with sequential N-implantation. Annealing temperature and annealing time were 1300°C and 60 min, respectively. Inset shows PL peak intensities corresponding to ABE as function of $[N]$.

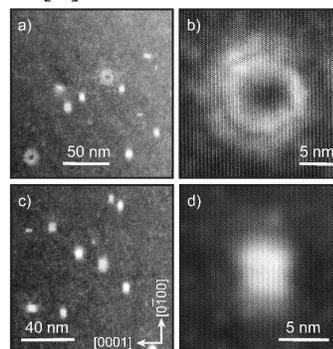


Fig. 10. LAADF-STEM images of (a) Mg-implanted GaN with sequential N-implantation, (c) N-implanted GaN ($[N]=1 \times 10^{18}$ cm^{-3}). Their magnified images are shown in (b) and (d), respectively. All samples were annealed at 1400°C ($t_a=5$ min).

IV. CONCLUSION

Annealing behaviors of vacancy-type defects in ion-implanted GaN were studied by positron annihilation. For as-implanted samples, the major defect species was identified as V_{Ga} -type defects. For N-implanted GaN, the size of the vacancies increased as the annealing temperature increased up to 1100°C and then shrank above 1200°C. This behavior was attributed to recombinations between V_{N} -type defects and implanted N atoms. For Mg-implanted GaN with sequential N-implantation, the major defect species after annealing above 1000°C was vacancy clusters such as $(V_{\text{Ga}}V_{\text{N}})_3$. Positively charged vacancy clusters were formed after annealing above 1000°C, and they had more V_{N} than neutral or negatively charged clusters. The effect of sequential N-implantation was most effective when the ratio of [N] to [Mg] was three. After annealing at 1400°C, the major secondary defects were nano-scale intrinsic defects and collapsed vacancy disks forming intrinsic dislocation loops, and the formation of the latter defects was suppressed in sequential N-implantation.

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