

1 **Performance Evaluation of a Diamond Quantum Magnetometer for Biomagnetic**  
2 **Sensing: A Phantom Study**

3 Naota Sekiguchi,<sup>1</sup> Yuta Kainuma,<sup>1</sup> Motofumi Fushimi,<sup>2</sup> Chikara Shinei,<sup>3,4</sup> Masashi  
4 Miyakawa,<sup>5</sup> Takashi Taniguchi,<sup>5</sup> Tokuyuki Teraji,<sup>3</sup> Hiroshi Abe,<sup>6</sup> Shinobu Onoda,<sup>6</sup>  
5 Takeshi Ohshima,<sup>6,7</sup> Mutsuko Hatano,<sup>1</sup> Masaki Sekino,<sup>2</sup> and Takayuki Iwasaki<sup>1</sup>

6 <sup>1</sup>*Department of Electrical and Electronic Engineering, Institute of Science Tokyo,*  
7 *Meguro, Tokyo 152-8550, Japan*

8 <sup>2</sup>*Department of Bioengineering, The University of Tokyo, Bunkyo, Tokyo 113-8656,*  
9 *Japan*

10 <sup>3</sup>*Research Center for Electronic and Optical Materials, National*  
11 *Institute for Materials Science, Tsukuba, Ibaraki 305-0044,*  
12 *Japan*

13 <sup>4</sup>*Department of Applied Physics, University of Tsukuba, Tsukuba, Ibaraki 305-8571,*  
14 *Japan*

15 <sup>5</sup>*Research Center for Materials Nanoarchitectonics, National*  
16 *Institute for Materials Science, Tsukuba, Ibaraki 305-0044,*  
17 *Japan*

18 <sup>6</sup>*Takasaki Institute for Advanced Quantum Science, National Institutes*  
19 *for Quantum Science and Technology, Takasaki, Gunma 370-1292,*  
20 *Japan*

21 <sup>7</sup>*Department of Materials Science, Tohoku University, Sendai, Miyagi 980-8579,*  
22 *Japan*

23 (\*sekiguchi.n.6ddd@m.isct.ac.jp)

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We employ a dry-type phantom to evaluate the performance of a diamond quantum magnetometer with a high sensitivity of about  $6 \text{ pT}/\sqrt{\text{Hz}}$  from the viewpoint of practical measurement in biomagnetic sensing. The dry phantom is supposed to represent an equivalent current dipole (ECD) generated by brain activity, emulating an encephalomagnetic field. The spatial resolution of the magnetometer is evaluated to be sufficiently higher than the length of the variation in the encephalomagnetic field distribution. The minimum detectable ECD moment is evaluated to be  $0.2 \text{ nA m}$  by averaging about 8000 measurements for a standoff distance of  $2.4 \text{ mm}$  from the ECD. We also discuss the feasibility of detecting an ECD in the measurement of an encephalomagnetic field in humans. We conclude that it is feasible to detect an encephalomagnetic field from a shallow cortex area such as the primary somatosensory cortex.

25 A diamond quantum magnetometer (DQM) based on nitrogen–vacancy (NV) center ensemble  
26 in diamond is a fascinating tool for biomagnetic sensing due to its favorable characteristics<sup>1,2</sup>.  
27 DQM can be operated at room temperature with a high sensitivity currently up to  $\text{pT}/\sqrt{\text{Hz}}$   
28 order<sup>3–5</sup>, which facilitates decreasing the distance to the measurement object. The short distance  
29 leads to a better spatial resolution of the target activity<sup>6</sup> and to a significantly larger signal because  
30 the biomagnetic field typically decays faster than the  $-1$  power of the distance<sup>7</sup>. The intrinsic  
31 spatial resolution of DQM itself is determined by the optically-excited volume of NV centers and  
32 can be decreased down to the sub-mm scale<sup>4,8,9</sup>. Additionally, a very wide dynamic range of  
33 DQM<sup>10–12</sup> provides the possibility of highly sensitive magnetometry in an ambient field outside a  
34 magnetic shield.

35 Sensitivity improvement in DQM has been actively studied<sup>13</sup> and actual application to biomag-  
36 netic sensing has been reported<sup>8,9,14</sup>, while few studies have reported on the evaluation of DQM  
37 from the viewpoint of biomagnetic sensing<sup>15</sup>. Many of those studies reported their field sensitiv-  
38 ity achieved<sup>3–5,16–22</sup>, but discussions about the stability<sup>4</sup> and the minimum detectable field in a  
39 biomagnetic sensing have generally been limited. The stability is of great importance because the  
40 current sensitivity around  $\sim \text{pT}/\sqrt{\text{Hz}}$  in DQM generally requires a long measurement time in total  
41 for accumulating a small signal. The minimum detectable field in this measurement depends on not  
42 only the measurement bandwidth and the number of accumulation but also characteristics of noise.  
43 Therefore, the evaluation of the minimum detectable field is essential to infer the performance of  
44 a magnetometer in practical applications. Furthermore, the evaluation of the intrinsic spatial reso-  
45 lution is important for a particular application including magnetoencephalography (MEG), where  
46 the estimation of the source generating a biomagnetic field by solving an inverse problem<sup>6,7</sup> would  
47 be disturbed if the intrinsic spatial resolution is worse than the length of the variation in the field  
48 distribution to be measured.

49 Here, we evaluate a DQM with a high sensitivity of about  $6 \text{ pT}/\sqrt{\text{Hz}}$  by using a dry type  
50 of phantom for an encephalomagnetic field in small animals. The dry phantom can emulate an  
51 encephalomagnetic field outside the brain and be considered as a representation of an equiva-  
52 lent current dipole (ECD) generated by brain activity. The spatial distribution of the phantom’s  
53 field was measured and agreed with the theoretical prediction, which indicates the intrinsic spatial  
54 resolution was sufficiently higher than the length of the variation in the encephalomagnetic field  
55 distribution. The minimum detectable field and the minimum detectable ECD in a typical mea-  
56 surement of stimulus-evoked field were investigated. It was found that the ECD moment of about

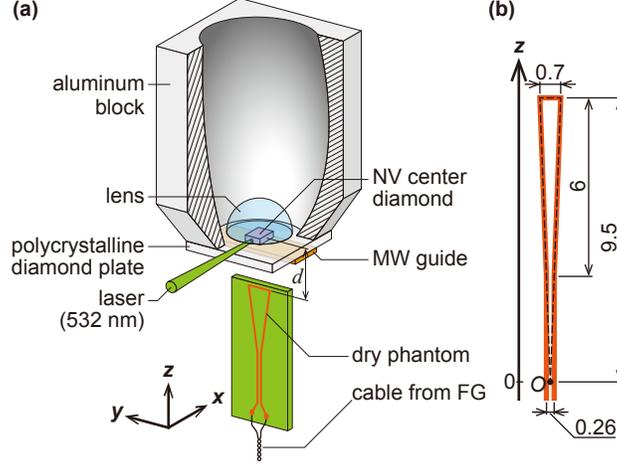


FIG. 1. (a) Experimental setup (not to scale) and (b) dimensions of the designed dry phantom.

57 0.5 nA m and of 0.2 nA m can be detected with the unity signal-to-noise ratio by averaging about  
 58 1500 and 8000 measurements, respectively, for a standoff distance of 2.4 mm. We also discuss the  
 59 feasibility of detecting an ECD in the measurement of an encephalomagnetic field in humans. We  
 60 conclude that it is feasible to detect a shallow ECD at, for example, the primary somatosensory  
 61 cortex area with our DQM.

62 A DQM used in this study is illustrated in Fig. 1(a). The DQM setup was almost the same  
 63 as that in our previous work<sup>4</sup> and inside a magnetically-shielded room. We used a single-crystal  
 64 (111) diamond synthesized by a high-pressure and high-temperature method. The initial con-  
 65 centration of substitutional nitrogen ( $N_s^0$ ) was controlled by using a titanium additive to a metal  
 66 solvent as a nitrogen getter<sup>23</sup>. NV centers were fabricated by electron beam irradiation and high-  
 67 temperature annealing. An electron spin resonance measurement yielded  $[NV^-] = 1.2$  ppm and  
 68  $[N_s^0] = 2.3$  ppm. The isotope ratio of  $^{13}C$  in the diamond was reduced to about 500 ppm. The  
 69 dephasing time,  $T_2^*$ , of the  $NV^-$  was estimated to be approximately  $2 \mu s^4$ .

70 An ensemble of  $NV^-$  in the diamond was excited from a side face by a green laser at 532 nm  
 71 with the power of 0.39 W. The excitation laser beam was focused onto the diamond and had the  
 72 spot size of about  $70 \mu m$  in diameter. The laser path length in the diamond was estimated to be  
 73 1 mm. The intrinsic spatial resolution of this DQM was therefore estimated to be  $70 \mu m$  and 1 mm  
 74 along the  $y$  and  $x$  direction, respectively. The laser-induced fluorescence from NV centers was  
 75 collected with a hemispherical lens and an elliptically-shaped inner wall of an aluminum block and  
 76 then detected by a photodiode. The intensity noise in the fluorescence due to the excitation-laser  
 77 intensity noise was reduced by a balanced detection technique<sup>4</sup>. Heat due to the laser illumination

78 was dissipated by attaching a polycrystalline diamond plate to the diamond.

79 We used a dry-type phantom that emulates an encephalomagnetic field from the brain of a  
80 small animal. While modeling an actual magnetic field generated from a neuron is quite difficult  
81 due to the complicated currents around the neuron, an analytical formula derived by Sarvas fairly  
82 reproduces the encephalomagnetic field under the assumptions as follows<sup>24</sup>: a field source, which  
83 is supposed to be an ensemble of intracellular currents at neurons, can be approximated as a single  
84 ECD  $\mathbf{Q}$  at  $\mathbf{r}_0$  in the spherically symmetric conductor with its center at the origin of the coordinate.  
85 The Sarvas' formula computes a magnetic field  $\mathbf{B}$  at the position  $\mathbf{r}$  of a sensor as

$$\mathbf{B}(\mathbf{r}) = \frac{\mu_0}{4\pi F^2} [F\mathbf{Q} \times \mathbf{r}_0 - \{(\mathbf{Q} \times \mathbf{r}_0) \cdot \mathbf{r}\} \nabla F], \quad (1)$$

86 where

$$F = |\mathbf{r} - \mathbf{r}_0| (|\mathbf{r} - \mathbf{r}_0||\mathbf{r}| + |\mathbf{r}|^2 - \mathbf{r}_0 \cdot \mathbf{r}). \quad (2)$$

87 The dry phantom that consists of an isosceles-triangle current is known as a source generating a  
88 magnetic field obeying the Sarvas' formula<sup>25</sup> and can emulate an encephalomagnetic field<sup>26</sup>. Our  
89 dry phantom made on a PCB was placed below the DQM with the distance  $d$  from the excited NV  
90 ensemble as shown in Fig. 1(a). The dimensions of the dry phantom [Fig. 1(b)] were determined  
91 by considering the size of the head of a small animal such as a rat. We intended to realize an  
92 isosceles-triangle coil with the base length  $l$  of 0.7 mm and the leg length of 9.5 mm, while the  
93 actual legs were connected to parallel wires at 6-mm away from the base. The ECD is supposed  
94 to be generated at the base with the moment  $Q_y$  along the  $y$  direction,  $Q_y = i_{\text{DP}}l$ , where  $i_{\text{DP}}$  is the  
95 current flowing on the phantom. The dry phantom was mounted on a  $z$  stage to vary the distance  
96  $d$  and on a motorized  $xy$  stage for the two-dimensional scan. We applied a sinusoidal current at  
97 33.33 Hz to this dry phantom to generate a test field. This frequency is within the primary band of  
98 an encephalomagnetic field<sup>27,28</sup>.

99 We performed a continuous-wave optically-detected magnetic resonance by applying a mi-  
100 crowave (MW) current through a MW guide on the other side of the polycrystalline diamond  
101 plate. A bias magnetic field of about 1 mT along the  $z$  axis was applied to the NV center ensem-  
102 ble by a permanent ring magnet. Here, the three possible resonances, which are associated with  
103 the hyperfine manifold, between the electron ground states  $|0\rangle \leftrightarrow |1\rangle$  were simultaneously driven  
104 by using a three-tone MW field<sup>8</sup>. The frequency of the three-tone field was modulated to employ  
105 lock-in detection. During the measurement of a magnetic field, we stabilized the MW frequency to  
106 the resonance frequency by monitoring the lock-in signal  $S_{\text{LI}}$  and applying a slow (2 Hz) PID servo

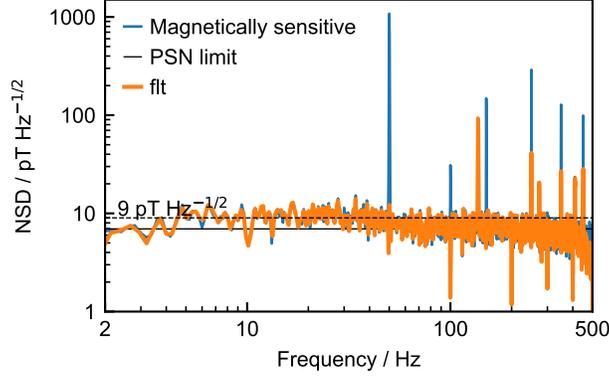


FIG. 2. Single-sided noise spectral density. The dashed line indicates the noise density floor of  $9 \text{ pT Hz}^{-1/2}$  at 33.33 Hz.

107 to an MW generator. The strength of the PID servo was optimized not to cause an increase in the  
 108 noise around the feedback frequency and to keep the resonance condition for a long period of mea-  
 109 surement. A magnetic field  $B_m$  faster than the PID servo was given as  $B_m = S_{\text{LI}} / (dS_{\text{LI}}/df \times \gamma_e)$ ,  
 110 where  $dS_{\text{LI}}/df$  is the zero-crossing slope around the resonance and  $\gamma_e = 28 \text{ GHz/T}$  is the gy-  
 111 romagnetic ratio. The minimum distance between the excited NV ensemble and a measurement  
 112 object was estimated to be 0.8 mm, limited by the thicknesses of the NV center diamond, the heat  
 113 spreading plate, and the MW guide.

114 The single-sided noise spectral density in the DQM was obtained from the 5-s measurement  
 115 with the cut-off frequency of 500-Hz at the lock-in detection as shown in Fig. 2. The noise spectral  
 116 density determined by photon shot-noise was calculated to be  $6.9 \text{ pT Hz}^{-1/2}$  (black solid line). The  
 117 measured noise density was found to be almost limited by the photon shot noise. We also found  
 118 that no noise peaks other than the power-line noises at 50 Hz and its harmonics were observed  
 119 at the encephalomagnetic-field frequency band  $< 100 \text{ Hz}$ <sup>27,28</sup>. The noise density floor around  
 120 the test-field frequency of 33.33 Hz was estimated to be  $9 \text{ pT Hz}^{-1/2}$  (dashed line). This noise  
 121 density floor corresponds to the field sensitivity of  $9 \text{ pT Hz}^{-1/2} / \sqrt{2} = 6 \text{ pT Hz}^{-1/2}$  on the widely  
 122 used definition of the sensitivity as  $\delta B \sqrt{T}$ <sup>8,19,29</sup>, where  $\delta B$  is the minimum detectable field for the  
 123 measurement time  $T$ .

124 We mapped the  $z$  component of the test field generated by the dry phantom by scanning the  $x$   
 125 and  $y$  positions of the phantom at five different distances as shown in Fig. 3. The heatmaps show  
 126 the amplitude of the sinusoidal test field obtained by least square method for the current peak  
 127 amplitude of 7.19 mA. At some positions (the hatched points), the measurement failed because the

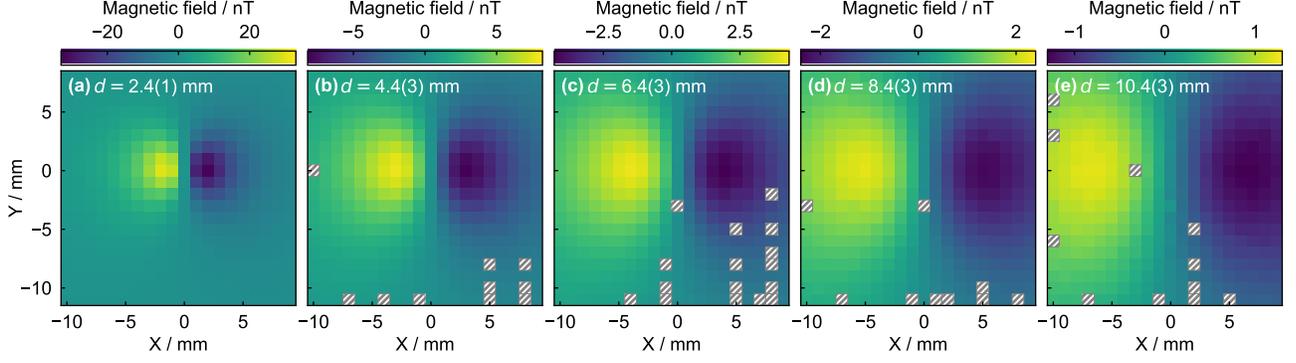


FIG. 3. Mapped magnetic field along the  $z$  axis with the different distance  $d$ . The measurement at the hatched regions failed due to a large noise from the motorized stage. The center of the phantom's base was aligned to the diamond center at  $(x, y) = (0, 0)$  by the eye.

128 motorized stage generated a magnetic field during the scanning that was larger than the applicable  
 129 range of the PID servo for the MW frequency. We observed two peaks with opposite polarity and  
 130 the sharper distribution for the closer distance, as expected from the ECD. The distance between  
 131 the two peaks approximately followed the distance  $d$ . The peak amplitude dependence on  $d$  shows  
 132 the decay faster than  $d^{-1}$ .

133 The field profiles along  $x$  axis at  $y = 0$  were extracted to compare the measured field with the nu-  
 134 merically calculated field by the Biot-Savart law and the Sarvas' formula [Eq. (1)]. Figure 4(a)-(e)  
 135 corresponds to the profile for Fig. 3(a)-(e), respectively. We found that the numerically calculated  
 136 field (black solid line) by the Biot-Savart law with the current of 7.19 mA along the designed  
 137 phantom's path agreed with the measured data (blue points) at all distances. This suggests that the  
 138 intrinsic spatial resolution of about 1 mm along the  $x$  direction was sufficient to resolve the peaks  
 139 of the phantom's field at  $d = 2.4 \pm 0.1$  mm, which is comparable to or shorter than the typical  
 140 standoff distance from the cortex of a small animal such as a rat<sup>30</sup>. The discrepancy between the  
 141 numerically calculated field and the field given by the Sarvas's formula (orange dashed line) at  
 142 larger distances indicates imperfection in our dry phantom. It was found by numerical calculation  
 143 that the discrepancy was attributed to two factors: the absence of the V-shaped wire compared to  
 144 the ideal isosceles triangle; and the presence of the slant wires from the parallel wire to pads for  
 145 connecting a cable [see Fig. 1(a)].

146 We investigated the minimum detectable magnetic field  $B_{\text{md}}$  in a typical measurement of  
 147 stimulus-evoked encephalomagnetic field. The stimulus-evoked field has a characteristic response

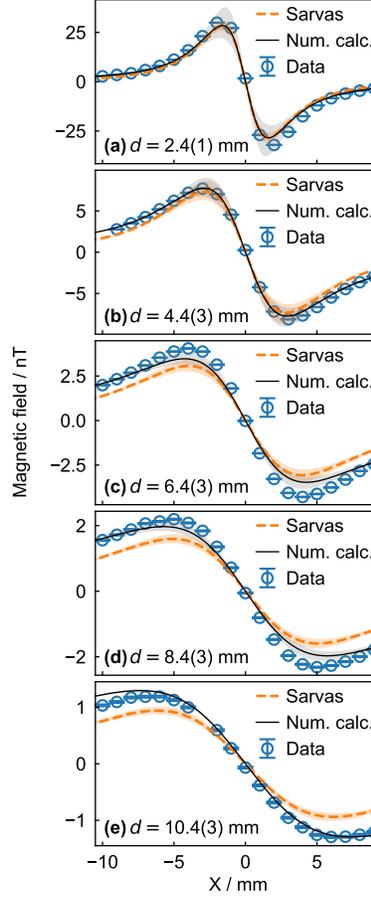


FIG. 4. Magnetic field profile along the  $x$  axis at  $y = 0$  for the different distance  $d$ . The measured field is shown by the blue circles. The black solid and orange dashed lines represent the numerically calculated profiles respectively by the Biot-Savart law and by the Sarvas' formula. The colored bands about the lines indicate the uncertainty caused by the uncertainty in  $d$ .

148 to a given stimulus, that is, its transient response to the stimulus is reproduced when the same  
 149 stimulus is applied. Typically, measurement of the evoked-field is repeated many times in syn-  
 150 chronization with the applied stimulus and averaged. In this evaluation of  $B_{md}$ , we placed the  
 151 phantom at  $(x, y) = (2.0, 0.0)$  mm with  $d = 2.4$  mm and repeatedly acquired the time trace for a  
 152 period of 570 ms. The total number of measurements was 8000 times, corresponds to the measure-  
 153 ment time of 76 minutes. The cut-off frequency of the low-pass filter at the lock-in amplifier was  
 154 set to 100 Hz, which corresponds to the frequency band of an encephalomagnetic field<sup>27,28</sup>. The  
 155 sinusoidal test current at 33.33 Hz was sent at the time  $t = 200$  ms for 8 periods to the phantom.  
 156 The peak amplitude of the current was  $i_{DP} = 0.69 \mu A$ . The corresponding moment of an ECD  
 157 was roughly estimated to be  $Q_y = 0.5$  nA m. It is supposed that small animals like rats generate

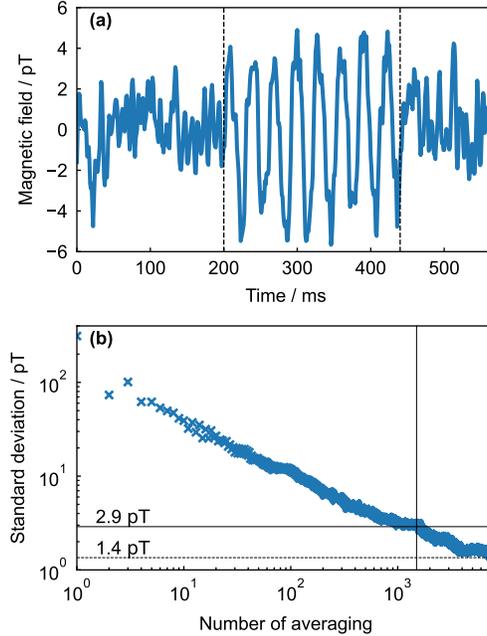


FIG. 5. Time domain measurement at  $(x, y) = (2.0, 0.0)$  mm. (a) Time trace averaged over 8000 measurements. The test field was generated at  $t = 200$  ms for 8 periods (between the dashed lines). (b) The standard deviation at  $t < 200$  ms as a function of the number of averaging.

158 such a small current dipole in the brain cortex by stimulation to, for example, auditory<sup>31</sup>. A type  
 159 of spontaneous brain activity including epilepsy should produce a stronger ECD<sup>32,33</sup>, but it is  
 160 difficult to average the field generated.

161 The time trace averaged over the 8000 acquisitions is shown in Fig. 5(a). The phantom's test  
 162 field was clearly observed and measured to have the root-mean-square amplitude of 2.9 pT, while  
 163 the standard deviation ( $B_{\text{md}}$ ) at  $t < 200$  ms was measured to be 1.4 pT. Furthermore, the minimum  
 164 detectable ECD moment for the 8000 times averaging and  $d = 2.4$  mm can be estimated to be  
 165  $0.5 \text{ nA m}/2.1 \simeq 0.2 \text{ nA m}$ .

166 We analyzed the decrease in the standard deviation in an averaged trace at  $t < 200$  ms as the  
 167 number of averaging was increased, as shown in Fig. 5(b). The 50 Hz power line noise quickly  
 168 decreased and was negligible comparing to the white noise at the number of averaging  $> 10$ . It  
 169 was found that the standard deviation scales as the number of averaging to the power of  $-1/2$ ,  
 170 which was enabled by the good noise performance. Since the standard deviation reached 2.9 pT at  
 171 around 1500 times averaging, we consider that our DQM can detect an ECD with  $Q_y \sim 0.5 \text{ nA m}$   
 172 by 15-minute measurement.

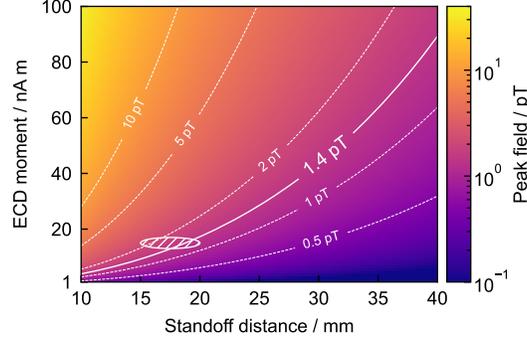


FIG. 6. Simulated peak field as a function of the standoff distance and the ECD moment. The white lines are the contour lines of the peak field. The hatched circle represents the previously reported ECD moment and standoff distance<sup>34</sup>.

173 Furthermore, we investigated the detectable ECD moment based on the Sarvas' formula for the  
 174 case of encephalomagnetic-field measurement of a human. Here, it was assumed for simplicity  
 175 that the ECD was located on the  $z$  axis at  $z_0$  and oriented along the tangential direction ( $x$ - $y$  plane).  
 176 Considering the size of a human brain, the distance between the center of the head-model sphere  
 177 and the DQM was fixed to 100 mm in this numerical calculation. The  $z$  component of the field was  
 178 numerically calculated by Eq. (1) and mapped by scanning the DQM position in the  $x$ - $y$  plane. We  
 179 explored the field maximum  $B_{z,\max}$  in the calculated field distribution for a given ECD moment  
 180 strength. Figure 6 shows  $B_{z,\max}$  as a function of the standoff distance and the ECD moment.  
 181 The standoff distance may be limited by the depth of the ECD from the head surface, since the  
 182 measurement distance of the DQM can be decreased to about 1 mm and negligible compared  
 183 to the depth. The minimum detectable field  $B_{\text{md}} = 1.4$  pT for the case of 8000 averaging is  
 184 indicated by the white solid contour line. The other contours in dashed line are guides for the eye.  
 185 It is supposed that a typical moment of an ECD in MEG measurement would be on the order of  
 186  $10 \text{ nA m}^7$ . The standoff distance largely varies on the activated region where the ECD is generated.  
 187 For example, for somatosensory stimulation, an ECD with the moment of about 15 nA m at the  
 188 depth of 15-20 mm (hatched circle) has been reported<sup>34</sup>. This simulation showed that the peak  
 189 field of the encephalomagnetic field that is generated by the previously reported ECD is stronger  
 190 than the minimum detectable field of 1.4 pT. Therefore, it is feasible for our DQM to detect an  
 191 encephalomagnetic field from such a shallow ECDs in humans.

192 We performed the measurement of a magnetic field generated by a dry-type phantom that em-  
 193 ulated an encephalomagnetic field in a small animal to evaluate a highly sensitive DQM from the

viewpoint of biomagnetic sensing. The single-sided noise spectral density of the DQM showed the very low noise floor of  $9 \text{ pT Hz}^{-1/2}$ , which corresponds to the sensitivity of about  $6 \text{ pT Hz}^{-1/2}$ . The spatial distribution of the phantom's field was measured by scanning the phantom relative to the DQM. The intrinsic spatial resolution of about 1 mm along the  $x$  direction of the DQM enabled us to observe the clear peaks of the phantom's field without smearing. For the case of time domain measurement, the minimum detectable field was found to be 1.4 pT with about 100-Hz bandwidth by averaging signal 8000 times. It was found that the ECD moment of about 0.2 nA m can be detected by averaging about 8000 measurements at the standoff distance of 2.4 mm. We also considered that it is feasible to detect an encephalomagnetic field from a shallow region of a human brain like the primary somatosensory cortex area by using our DQM. Although the estimations of the position and the moment of an ECD were not performed due to the discrepancy of our dry phantom from an ideal one, it is demonstrated that the evaluation of a DQM using a phantom is beneficial from the viewpoint of biomagnetic sensing.

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## REFERENCES

- <sup>1</sup>T. Zhang, G. Pramanik, K. Zhang, M. Gulka, L. Wang, J. Jing, F. Xu, Z. Li, Q. Wei, P. Cigler, and Z. Chu, "Toward Quantitative Bio-sensing with Nitrogen-Vacancy Center in Diamond," *ACS Sensors* **6**, 2077–2107 (2021).
- <sup>2</sup>N. Aslam, H. Zhou, E. K. Urbach, M. J. Turner, R. L. Walsworth, M. D. Lukin, and H. Park, "Quantum sensors for biomedical applications," *Nature Reviews Physics* **5**, 157–169 (2023).
- <sup>3</sup>J. F. Barry, M. H. Steinecker, S. T. Alsid, J. Majumder, L. M. Pham, M. F. O'Keeffe, and D. A. Braje, "Sensitive ac and dc magnetometry with nitrogen-vacancy-center ensembles in diamond," *Physical Review Applied* **22**, 044069 (2024).
- <sup>4</sup>N. Sekiguchi, M. Fushimi, A. Yoshimura, C. Shinei, M. Miyakawa, T. Taniguchi, T. Teraji, H. Abe, S. Onoda, T. Ohshima, M. Hatano, M. Sekino, and T. Iwasaki, "Diamond quantum magnetometer with dc sensitivity of sub-10 pT Hz<sup>-1/2</sup> toward measurement of biomagnetic field," *Physical Review Applied* **21**, 064010 (2024).
- <sup>5</sup>H. Wang, K. L. Tiwari, K. Jacobs, M. Judy, X. Zhang, D. R. Englund, and M. E. Trusheim, "A spin-refrigerated cavity quantum electrodynamic sensor," *Nature Communications* **15**, 10320

- 224 (2024).
- 225 <sup>6</sup>R. Körber, J.-H. Storm, H. Seton, J. P. Mäkelä, R. Paetau, L. Parkkonen, C. Pfeiffer, B. Riaz,  
226 J. F. Schneiderman, H. Dong, S.-m. Hwang, L. You, B. Inglis, J. Clarke, M. A. Espy, R. J.  
227 Ilmoniemi, P. E. Magnelind, A. N. Matlashov, J. O. Nieminen, P. L. Volegov, K. C. J. Zeven-  
228 hoven, N. Höfner, M. Burghoff, K. Enpuku, S. Y. Yang, J.-J. Chieh, J. Knuutila, P. Laine, and  
229 J. Nenonen, “SQUIDS in biomagnetism: A roadmap towards improved healthcare,” *Supercon-*  
230 *ductor Science and Technology* **29**, 113001 (2016).
- 231 <sup>7</sup>M. Hämäläinen, R. Hari, R. J. Ilmoniemi, J. Knuutila, and O. V. Lounasmaa,  
232 “Magnetoencephalography—theory, instrumentation, and applications to noninvasive studies of  
233 the working human brain,” *Reviews of Modern Physics* **65**, 413–497 (1993).
- 234 <sup>8</sup>J. F. Barry, M. J. Turner, J. M. Schloss, D. R. Glenn, Y. Song, M. D. Lukin, H. Park, and R. L.  
235 Walsworth, “Optical magnetic detection of single-neuron action potentials using quantum de-  
236 fects in diamond,” *Proceedings of the National Academy of Sciences* **113**, 14133–14138 (2016).
- 237 <sup>9</sup>K. Arai, A. Kuwahata, D. Nishitani, I. Fujisaki, R. Matsuki, Y. Nishio, Z. Xin, X. Cao, Y. Hatano,  
238 S. Onoda, C. Shinei, M. Miyakawa, T. Taniguchi, M. Yamazaki, T. Teraji, T. Ohshima,  
239 M. Hatano, M. Sekino, and T. Iwasaki, “Millimetre-scale magnetocardiography of living rats  
240 with thoracotomy,” *Communications Physics* **5**, 200 (2022).
- 241 <sup>10</sup>H. Clevenson, L. M. Pham, C. Teale, K. Johnson, D. Englund, and D. Braje, “Robust  
242 high-dynamic-range vector magnetometry with nitrogen-vacancy centers in diamond,” *Applied*  
243 *Physics Letters* **112**, 252406 (2018).
- 244 <sup>11</sup>Y. Hatano, J. Shin, J. Tanigawa, Y. Shigenobu, A. Nakazono, T. Sekiguchi, S. Onoda,  
245 T. Ohshima, K. Arai, T. Iwasaki, and M. Hatano, “High-precision robust monitoring of  
246 charge/discharge current over a wide dynamic range for electric vehicle batteries using diamond  
247 quantum sensors,” *Scientific Reports* **12**, 13991 (2022).
- 248 <sup>12</sup>C. Wang, Q. Liu, Y. Hu, F. Xie, K. Krishna, N. Wang, L. Wang, Y. Wang, K. C. Toussaint,  
249 J. Cheng, H. Chen, and Z. Wu, “Realization of high-dynamic-range broadband magnetic-field  
250 sensing with ensemble nitrogen-vacancy centers in diamond,” *Review of Scientific Instruments*  
251 **94**, 015109 (2023).
- 252 <sup>13</sup>J. F. Barry, J. M. Schloss, E. Bauch, M. J. Turner, C. A. Hart, L. M. Pham, and R. L. Walsworth,  
253 “Sensitivity optimization for NV-diamond magnetometry,” *Reviews of Modern Physics* **92**,  
254 015004 (2020).
- 255 <sup>14</sup>Z. Yu, Y. Xie, G. Jin, Y. Zhu, Q. Zhang, F. Shi, F.-y. Wan, H. Luo, A.-h. Tang, and X. Rong,

- 256 “Noninvasive magnetocardiography of a living rat based on a diamond quantum sensor,” *Physical Review Applied* **21**, 064028 (2024).
- 257
- 258 <sup>15</sup>C. Zhang, J. Zhang, M. Widmann, M. Benke, M. Kübler, D. Dasari, T. Klotz, L. Gizzi, O. Röhrle,  
259 P. Brenner, and J. Wrachtrup, “Optimizing NV magnetometry for Magnetoneurography and  
260 Magnetomyography applications,” *Frontiers in Neuroscience* **16**, 1034391 (2023).
- 261 <sup>16</sup>I. Fescenko, A. Jarmola, I. Savukov, P. Kehayias, J. Smits, J. Damron, N. Ristoff, N. Mosavian,  
262 and V. M. Acosta, “Diamond magnetometer enhanced by ferrite flux concentrators,” *Physical*  
263 *Review Research* **2**, 023394 (2020).
- 264 <sup>17</sup>X. Gao, C. Yu, S. Zhang, H. Lin, J. Guo, M. Ma, Z. Feng, and F.-W. Sun, “High sensitivity of  
265 diamond nitrogen-vacancy magnetometer with magnetic flux concentrators via enhanced fluo-  
266 rescence collection,” *Diamond and Related Materials* **139**, 110348 (2023).
- 267 <sup>18</sup>S. Graham, A. Rahman, L. Munn, R. Patel, A. Newman, C. Stephen, G. Colston, A. Nikitin,  
268 A. Edmonds, D. Twitchen, M. Markham, and G. Morley, “Fiber-Coupled Diamond Magnetom-  
269 etry with an Unshielded Sensitivity of 30 pT / Hz,” *Physical Review Applied* **19**, 044042 (2023).
- 270 <sup>19</sup>J. M. Schloss, J. F. Barry, M. J. Turner, and R. L. Walsworth, “Simultaneous Broadband Vector  
271 Magnetometry Using Solid-State Spins,” *Physical Review Applied* **10**, 034044 (2018).
- 272 <sup>20</sup>T. Wolf, P. Neumann, K. Nakamura, H. Sumiya, T. Ohshima, J. Isoya, and J. Wrachtrup, “Sub-  
273 picotesla Diamond Magnetometry,” *Physical Review X* **5**, 041001 (2015).
- 274 <sup>21</sup>Y. Xie, H. Yu, Y. Zhu, X. Qin, X. Rong, C.-K. Duan, and J. Du, “A hybrid magnetometer towards  
275 femtotesla sensitivity under ambient conditions,” *Science Bulletin* **66**, 127–132 (2021).
- 276 <sup>22</sup>C. Zhang, F. Shagieva, M. Widmann, M. Kübler, V. Vorobyov, P. Kapitanova, E. Nenasheva,  
277 R. Corkill, O. Rhrle, K. Nakamura, H. Sumiya, S. Onoda, J. Isoya, and J. Wrachtrup, “Diamond  
278 Magnetometry and Gradiometry Towards Subpicotesla dc Field Measurement,” *Physical Review*  
279 *Applied* **15**, 064075 (2021).
- 280 <sup>23</sup>M. Miyakawa, C. Shinei, and T. Taniguchi, “Nitrogen concentration control in diamonds grown  
281 in Co–(Fe)–Ti/Al solvents under high-pressure and high-temperature,” *Japanese Journal of Ap-*  
282 *plied Physics* **61**, 045507 (2022).
- 283 <sup>24</sup>J. Sarvas, “Basic mathematical and electromagnetic concepts of the biomagnetic inverse prob-  
284 lem,” *Physics in Medicine and Biology* **32**, 11–22 (1987).
- 285 <sup>25</sup>G. Uehara, M. Higuchi, Y. Adachi, J. Kawai, H. Tanaka, Y. Haruta, M. Miyamoto, T. Shimozu,  
286 and M. Kawabata, “An error in magnetic field of isosceles-triangle coil as a representation of  
287 equivalent current dipole,” *International Congress Series* **1300**, 607–610 (2007).

- 288 <sup>26</sup>D. Oyama, Y. Adachi, M. Yumoto, I. Hashimoto, and G. Uehara, “Dry phantom for magnetoen-  
289 cephalography —Configuration, calibration, and contribution,” *Journal of Neuroscience Meth-*  
290 *ods* **251**, 24–36 (2015).
- 291 <sup>27</sup>R. Hari, S. Baillet, G. Barnes, R. Burgess, N. Forss, J. Gross, M. Hämäläinen, O. Jensen,  
292 R. Kakigi, F. Mauguière, N. Nakasato, A. Puce, G.-L. Romani, A. Schnitzler, and S. Taulu,  
293 “IFCN-endorsed practical guidelines for clinical magnetoencephalography (MEG),” *Clinical*  
294 *Neurophysiology* **129**, 1720–1747 (2018).
- 295 <sup>28</sup>P. J. Uhlhaas, T. Grent-’t-Jong, and J. Gross, “Magnetoencephalography and Translational Neu-  
296 roscience in Psychiatry,” *JAMA Psychiatry* **75**, 969 (2018).
- 297 <sup>29</sup>J. M. Taylor, P. Cappellaro, L. Childress, L. Jiang, D. Budker, P. R. Hemmer, A. Yacoby,  
298 R. Walsworth, and M. D. Lukin, “High-sensitivity diamond magnetometer with nanoscale reso-  
299 lution,” *Nature Physics* **4**, 810–816 (2008).
- 300 <sup>30</sup>G. Paxinos and C. Watson, *The Rat Brain in Stereotaxic Coordinates*, 4th ed. (Academic Press,  
301 1998).
- 302 <sup>31</sup>J.-E. Kim, I.-S. Kim, K. Kim, S. Lim, H. Kwon, C. S. Kang, S. Ahn, K. K. Yu, and Y.-H.  
303 Lee, “Development of a bio-magnetic measurement system and sensor configuration analysis  
304 for rats,” *Review of Scientific Instruments* **88**, 044704 (2017).
- 305 <sup>32</sup>Q. Chen, M. Chopp, H. Chen, and N. Tepley, “Magnetoencephalography of focal cerebral is-  
306 chemia in rats.” *Stroke* **23**, 1299–1303 (1992).
- 307 <sup>33</sup>O. Alem, A. M. Benison, D. S. Barth, J. Kitching, and S. Knappe, “Magnetoencephalogra-  
308 phy of Epilepsy with a Microfabricated Atomic Magnetode,” *The Journal of Neuroscience* **34**,  
309 14324–14327 (2014).
- 310 <sup>34</sup>R. Körber, J. O. Nieminen, N. Höfner, V. Jazbinšek, H.-J. Scheer, K. Kim, and M. Burghoff, “An  
311 advanced phantom study assessing the feasibility of neuronal current imaging by ultra-low-field  
312 NMR,” *Journal of Magnetic Resonance* **237**, 182–190 (2013).