

Thermoelectric Properties of $\text{Bi}_{1-x}\text{Pb}_x\text{Cu}_{1-x}\text{SeO}$ Oxyselenides[†]

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In this work, $\text{Bi}_{1-x}\text{Pb}_x\text{Cu}_{1-x}\text{SeO}$ ($x = 0, 0.02, 0.06, \text{ and } 0.08$) were synthesized by solid-state reaction followed by spark plasma sintering. The effect of simultaneous Bi to Pb substitution and Cu vacancies introduction on thermoelectric properties was investigated systematically. The power factor was significantly enhanced, contributing to the increase in the zT value. As a result, the zT_{max} of 0.75 at 773 K was obtained for $\text{Bi}_{0.94}\text{Pb}_{0.06}\text{Cu}_{0.94}\text{SeO}$ sample. To reveal the factors constraining the zT_{max} in $\text{Bi}_{1-x}\text{Pb}_x\text{Cu}_{1-x}\text{SeO}$ -based oxyselenides, a careful analysis of literature data was performed. We highlighted that in Pb-doped oxyselenides, the power factor is almost independent of the synthesis technique, while the lattice thermal conductivity is the main property determining zT_{max} and is highly affected by the synthesis method.

Introduction

Over the past few decades, global energy consumption has dramatically increased, with about two-thirds of produced energy being lost as waste heat. Therefore, technologies capable of efficiently recuperating this heat are crucial to achieving sustainable development goals. In this context, thermoelectric materials, which enable direct conversion of waste heat into electricity, have attracted increasing attention from the scientific community.¹ The efficiency of thermoelectric energy conversion depends on the material's transport properties and is determined by the thermoelectric figure of merit $zT = \alpha^2 \sigma T / \kappa_{\text{tot}}$, where α , σ , T , and κ_{tot} represent the Seebeck coefficient, electrical conductivity, absolute temperature, and total thermal conductivity, respectively.²

Currently, p -type BiCuSeO-based oxyselenides are among the most efficient oxygen-containing thermoelectric materials.³ BiCuSeO crystallizes in a tetragonal layered ZrCuSiAs structure type with a $P4/nmm$ space group and two formula units per unit cell. The crystal structure consists of alternately stacked insulating $(\text{Bi}_2\text{O}_2)^{2+}$ layers (so-called "charge reservoir" layers) and conducting $(\text{Cu}_2\text{Se}_2)^{2-}$ layers along the c -axis.^{4,5} These BiCuSeO-based oxyselenides have drawn interest as thermoelectric materials, primarily due to their intrinsically low thermal conductivity, not exceeding $1.5 \text{ W m}^{-1} \text{ K}^{-1}$ at room temperature, coupled with a relatively high Seebeck coefficient.⁶ The main focus in enhancing zT of BiCuSeO is optimizing the charge carrier concentration n , which significantly boosts the power factor ($\alpha^2 \sigma$) by balancing α and σ values. Employing this approach in BiCuSeO-based oxyselenides has yielded $zT_{\text{max}} > 1.2$ at $T \geq 773 \text{ K}$ achieved through the substitution of bismuth with various elements.^{7,8}

However, as identified by Ioffe over six decades ago, the efficient thermoelectric material requires not only an optimal charge concentration n_{opt} but also a maximized ratio of charge carrier mobility μ to lattice thermal conductivity κ_{lat} , as this determined the zT_{max} achievable at a given temperature and charge carrier

concentration.² While optimizing charge carrier concentration in BiCuSeO is feasible through heterovalent Bi substitution, methods to maximize $\mu / \kappa_{\text{lat}}$ are less straightforward. One of the peculiar approaches was recently indicated to be beneficial for keeping a high n/μ ratio by promoting the carriers' delocalization between the charge reservoir and conducting layers via introducing Bi and/or Cu vacancies. Although interlayer delocalization of carriers is favorable for interlayer charge transfer, conventional Bi substitution is still necessary to provide a sufficient number of charge carriers for diffusion.^{9,10} In this regard, Pb has been identified as the most effective dopant for oxyselenides in terms of the number of charge carriers per one Pb atom introduced into the system during doping.¹¹ Concurrently, copper vacancies are the predominant defect in BiCuSeO responsible for its native p -type self-doping.^{12,13}

In our study, we synthesized a series of samples with the nominal composition of $\text{Bi}_{1-x}\text{Pb}_x\text{Cu}_{1-x}\text{SeO}$ ($x = 0, 0.02, 0.06, \text{ and } 0.08$) to achieve simultaneous optimization of charge carrier concentration while maintaining high values of their mobility by combining the substitution of bismuth with lead (Pb_{Bi}) and introducing extra copper vacancies (V_{Cu}). $\text{Bi}_{1-x}\text{Pb}_x\text{Cu}_{1-x}\text{SeO}$ -based oxyselenides are among the most studied, and the influence of Bi to Pb substitution on transport properties has been thoroughly investigated. We demonstrated that the key factor affecting the thermoelectric performance of Pb-doped BiCuSeO and the variability reported in zT_{max} values is, in fact, the $\mu / \kappa_{\text{lat}}$ ratio. At the same time, the power factor is predominantly influenced by the Pb_{Bi} concentration and is weakly dependent on the concentration of vacancies and/or the second dopant. Ultimately, we observed an almost fourfold increase in the power factor for $\text{Bi}_{1-x}\text{Pb}_x\text{Cu}_{1-x}\text{SeO}$ with $x = 0.06$ and 0.08 at 773 K. The $\mu_w / \kappa_{\text{lat}}$ ratio doubled with an increase from $x = 0$ to $x = 0.02$, yet further increase in x did not yield an increase in $\mu_w / \kappa_{\text{lat}}$ mainly due to the increase in κ_{lat} for $x > 0.02$. Ultimately, a notable increase zT by a factor of three was realized, and a $zT_{\text{max}} = 0.75$ at 773 K for $\text{Bi}_{0.94}\text{Pb}_{0.06}\text{Cu}_{0.94}\text{SeO}$ was achieved.

Materials and Methods

Compounds with the nominal chemical compositions of $\text{Bi}_{1-x}\text{Pb}_x\text{Cu}_{1-x}\text{SeO}$ ($x = 0, 0.02, 0.06, \text{ and } 0.08$) were synthesized as reported elsewhere¹⁴ using high purity CuO, Bi, Se, and PbO powders as the raw materials. The obtained products were consolidated through spark plasma sintering (SPS) at 903 K for 5 min under an axial pressure of 50 MPa. Phase composition analysis

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was carried out through powder X-ray diffraction (PXRD) using a DRON-3 diffractometer (IC Bouvestnik, Russia) with a $\text{Cu-K}\alpha$ radiation ($\lambda = 1.54178 \text{ \AA}$). The PXRD patterns were refined by the Rietveld method with the self-developed software package.¹⁵ Scanning electron microscopy (SEM) and energy-dispersive X-ray spectroscopy (EDS) were performed using a Vega 3SB SEM (Tescan, Czech Republic) equipped with an x-act EDS detector (Oxford Instruments, UK). The thermal diffusivity χ was measured in the axial direction of disk-shaped samples of $\varnothing 10 \times 1 \text{ mm}^2$ using a laser flash analyzer (LFA 467 Hyperflash, Netzsch, Germany). The total thermal conductivity κ_{tot} was calculated using the formula $\kappa_{tot} = \chi C_p d$, where C_p represents the specific heat capacity which was estimated using the comparison method (with pyroceram-9060 as a reference sample), and is in a good agreement with the C_p value calculated by the Debye model. The density d was determined using the Archimedes method. The Seebeck coefficient α and the electrical conductivity σ were measured in the radial direction of bar-shaped specimens with dimensions of $10 \times 3 \times 1 \text{ mm}^3$ using a commercial apparatus (ZEM-3, Advance Riko Inc., Japan) under a partial He atmosphere from 300 to 773 K. The uncertainty in σ and α was 5% and 4%, respectively, along with 8% in κ_{tot} . The overall uncertainty in zT was estimated to be less than 17%. The error bars are not shown for clarity and to enhance the visual presentation of the figures.

Results and Discussion

The PXRD patterns of $\text{Bi}_{1-x}\text{Pb}_x\text{Cu}_{1-x}\text{SeO}$ ($x = 0, 0.02, 0.06, \text{ and } 0.08$) samples are presented in Fig. 1a. All major reflections can be indexed to the tetragonal BiCuSeO phase (PDF#01-076-6689) with a ZrCuSiAs structure type and $P4/mmm$ space group. A minor amount of the Bi_2O_3 (PDF#01-088-2043) secondary phase ($\leq 4.5 \text{ vol.}\%$) was detected in all the samples, indicating significant Cu deficiency as expected.^{12,13,16,17} In the case of dual doping, the evolution of lattice parameters can be affected by two competing factors. On the one hand, Cu deficiency tends to decrease both lattice parameters a and c ,¹⁷ while on the other hand, a larger ionic radius of Pb^{2+} (1.19 \AA) as compared to that of Bi^{3+} (1.03 \AA)¹⁸ should result in an increase of lattice parameters.^{19,20} Our experimental results reveal that, even with a simultaneous increase in the concentration of Pb_{Bi} and V_{Cu} , both lattice parameters a and c exhibit a gradual increase (Fig. 1b), resulting in the expansion of the unit cell volume. However, it is noteworthy that the effect of copper vacancies is noticeable, and for samples with the same nominal concentration x , the lattice parameters of dually doped specimens are slightly smaller than those of single Pb-doped ones.¹⁹⁻²¹

Figure 2 displays an SEM micrograph, providing an overview of the microstructure representative of all samples. Randomly arranged platelet grains are stacked densely, which is reflected by the high relative densities for all samples ($\geq 95\%$ of theoretical density). The thickness of the grains ranges from 300 to 600 nm. EDS analysis confirms that the actual compositions of the samples closely align with the nominal ones (Table S1).

Figure 3a shows the temperature dependence of electrical conductivity σ of $\text{Bi}_{1-x}\text{Pb}_x\text{Cu}_{1-x}\text{SeO}$ ($x = 0, 0.02, 0.06, \text{ and } 0.08$) samples. For all the samples, σ gradually decreases with

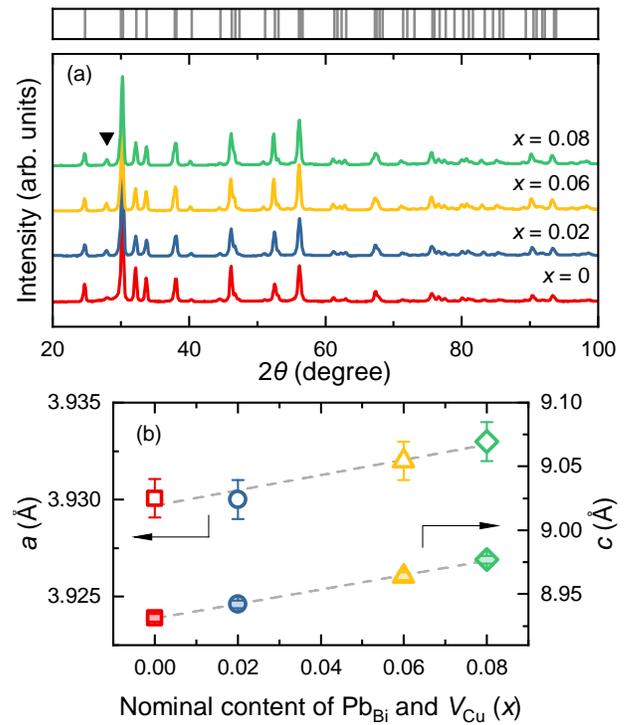


Figure 1 (a) Powder XRD patterns of the $\text{Bi}_{1-x}\text{Pb}_x\text{Cu}_{1-x}\text{SeO}$ ($x = 0, 0.02, 0.06, \text{ and } 0.08$) samples. Bragg's reflections for the BiCuSeO phase are indicated by gray ticks on the top part of the figure. Diffraction peak corresponding to the Bi_2O_3 secondary phase is indicated by a black solid triangle (\blacktriangledown). (b) Lattice parameters (a and c) as a function of nominal concentration of Pb_{Bi} and V_{Cu} .

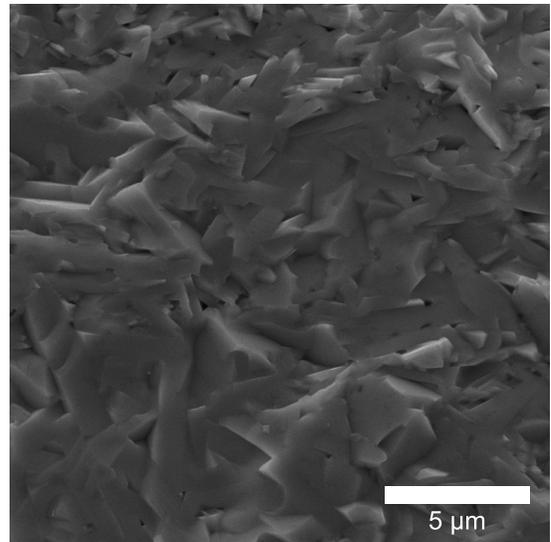


Figure 2 SEM micrograph of the fractured cross-section of $\text{Bi}_{0.94}\text{Pb}_{0.06}\text{Cu}_{0.94}\text{SeO}$ sample.

the increase of temperature, exhibiting a metal-like behavior as expected for heavily doped semiconductors.²² As the content of Pb_{Bi} and V_{Cu} increases from 0 to 0.08, σ is vastly improved from $\sim 30 \Omega^{-1} \text{ cm}^{-1}$ for the pristine BiCuSeO sample to $\sim 450 \Omega^{-1} \text{ cm}^{-1}$ for $\text{Bi}_{0.92}\text{Pb}_{0.08}\text{Cu}_{0.92}\text{SeO}$ at room temperature. Considering only a slight improvement in the weighted mobility

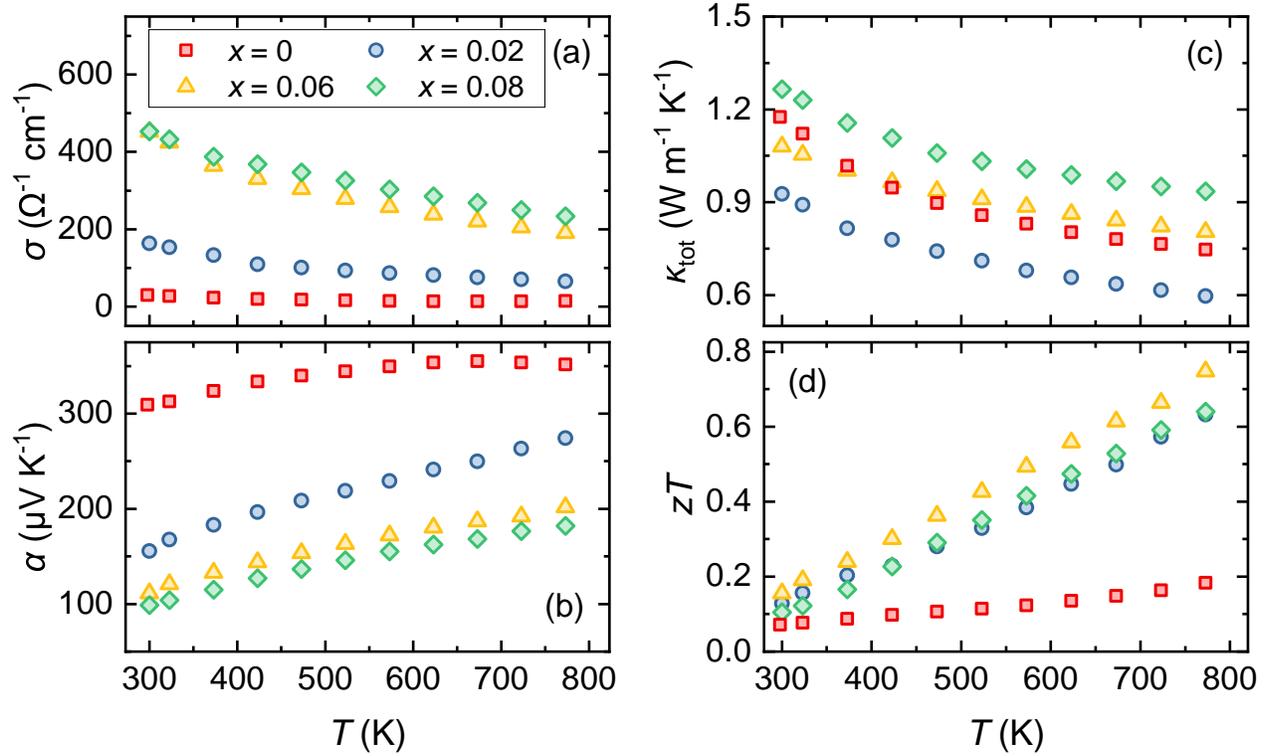


Figure 3 Temperature dependence of the (a) electrical conductivity σ , (b) Seebeck coefficient α , (c) total thermal conductivity κ_{tot} , and (d) the figure of merit zT for $\text{Bi}_{1-x}\text{Pb}_x\text{Cu}_{1-x}\text{SeO}$ ($x = 0, 0.02, 0.06,$ and 0.08) samples.

μ_w ($\mu_w \approx \mu (m_d^*/m_e)^{3/2}$, where m_d^* is the density of states effective mass, and m_e is the electron mass) as shown in Fig. S1, the main reason for such a significant enhancement of the electrical conductivity is from the increase in the charge carriers concentration, which is also in agreement with previous reports.^{21,23} Furthermore, the carriers density increase originates mostly from the substitution of Bi^{3+} with Pb^{2+} , as even introducing only dual vacancies V_{Bi} and V_{Cu} in BiCuSeO led to an order of magnitude lower conductivity in comparison with Pb -substituted ones.¹⁰ Nonetheless, as discussed in the Introduction section, the introduction of vacancies plays not a primary but a complementary role, leading to a more pronounced delocalization of charge carriers between the layers, thereby promoting interlayer charge transfer.⁹ This, in turn, allows for the retention of higher charge carrier mobility values at the same level of charge carrier concentration (see Fig. 12 in Ref. [11]).

Figure 3b illustrates the Seebeck coefficient α as a function of temperature for $\text{Bi}_{1-x}\text{Pb}_x\text{Cu}_{1-x}\text{SeO}$ ($x = 0, 0.02, 0.06,$ and 0.08) samples. The positive Seebeck coefficient for all samples in the entire temperature range confirms p -type conduction, indicating that holes are the majority charge carriers. The highest Seebeck coefficient throughout the investigated temperature range is observed in nominally undoped BiCuSeO . However, with the increase in Pb_{Bi} and V_{Cu} concentration, α gradually decreases consistently with the changes in σ (Fig. S2). Considering the relationship between the Seebeck coefficient and carrier concentration in semiconductors,² such a simultaneous change in electrical conductivity and the Seebeck coefficient further underscores

that it is primarily driven by a significant increase in charge carrier concentration, as was discussed previously. Consequently, at 773 K, α reduced from $352 \mu\text{V K}^{-1}$ for $x = 0$ to $182 \mu\text{V K}^{-1}$ for $x = 0.08$. The achieved optimal balance between σ and α results in a significant enhancement in the power factor ($\alpha^2\sigma$), showing more than a fourfold increase at 773 K, from $1.8 \mu\text{W cm}^{-1} \text{K}^{-2}$ to $7.7 \mu\text{W cm}^{-1} \text{K}^{-2}$, as the Pb_{Bi} and V_{Cu} content is raised from $x = 0$ to 0.08 (Fig. 4a, Fig. S3).

Figure 3c shows the temperature-dependent total thermal conductivity $\kappa_{tot}(T)$ of $\text{Bi}_{1-x}\text{Pb}_x\text{Cu}_{1-x}\text{SeO}$ ($x = 0, 0.02, 0.06,$ and 0.08) samples. For the pristine BiCuSeO , κ_{tot} decreases from $1.17 \text{ W m}^{-1} \text{K}^{-1}$ at 300 K to $0.75 \text{ W m}^{-1} \text{K}^{-1}$ at 773 K. Following simultaneous Bi for Pb substitution and Cu vacancies introduction, κ_{tot} initially decreases ($x = 0.02$) and then rises to $0.93 \text{ W m}^{-1} \text{K}^{-1}$ at 773 K ($x = 0.08$). This can be understood by considering evolution in both the lattice κ_{lat} and electronic κ_{el} components of total thermal conductivity. The κ_{lat} is calculated from $\kappa_{lat} = \kappa_{tot} - \kappa_{el}$, where κ_{el} can be determined through the Wiedemann–Franz law, expressed as $\kappa_{el} = L\sigma T$, with L representing the Lorentz number.² The Lorentz number is estimated with sufficient accuracy within the single parabolic band model with acoustic phonon scattering from experimental α values as $L = 1.5 + \exp\{-|\alpha|/116\}$.²⁴ Owing to a significant enhancement in σ , the changes in κ_{tot} become more apparent. In pristine BiCuSeO , κ_{el} contributes to approximately 2% of the total, but this contribution rises up to over 30% with an increase in Pb_{Bi} and V_{Cu} concentration (Fig. S4). Analysis of the $\kappa_{lat}(T)$ reveals that point defects are the dominant phonon scattering mechanism, ev-

identified by a decrease of κ_{lat} with temperature as $T^{-0.5}$ for all the samples (Fig. S5). However, for the samples with $x > 0.02$ κ_{lat} slightly increases but still remains lower than that of pristine BiCuSeO. This increase could be attributed to an increase in the grain size, sound velocities, and/or Pb–O bond strength due to the introduction of Pb_{Bi} and V_{Cu} .²⁵

Combining the electrical and thermal transport properties, the figure of merit zT as a function of temperature is calculated and displayed in Fig. 3d. For $Bi_{1-x}Pb_xCu_{1-x}SeO$ samples ($x > 0$), the zT values are enhanced over the entire test temperature range by more than 3 times. This improvement is mainly attributed to a significant boost in the power factor, leading to a zT_{max} value of 0.75 at 773 K for $Bi_{0.94}Pb_{0.06}Cu_{0.94}SeO$. On the one hand, a nearly threefold increase in zT was achieved compared to undoped BiCuSeO. On the other hand, we must admit that this is not the highest zT value achieved for $Bi_{1-x}Pb_xCu_{1-x}SeO$ -based oxy-selenides. Furthermore, a detailed analysis of literature data reveals that reported zT values for Pb-doped BiCuSeO are scattered between 0.56 and 1.05 at $T = 773$ K and $x = 0.06$. Considering that the power factor is practically independent of the synthesis method (microstructure: grain size and shape, defects, etc.), as shown in Fig. 4a, the main factor determining zT value is the lattice thermal conductivity (Fig. S6). After careful investigation, it can be concluded that the highest zT values were obtained in samples where lattice thermal conductivity was suppressed by about twice compared to the glassy limit (≈ 0.59 W m⁻¹ K⁻¹). Interestingly, grain size is not the determining factor in reducing κ_{lat} . The lowest κ_{lat} values were reported for samples synthesized by methods like mechanochemical synthesis,²⁶ self-propagating high-temperature synthesis,^{27,28} and high-energy ball milling,¹⁶ all of them leading to the formation of various defects, such as nanodots,¹⁶ amorphous regions,^{27,28} etc.

Conclusions

The thermoelectric properties of the $Bi_{1-x}Pb_xCu_{1-x}SeO$ ($x = 0, 0.02, 0.06, \text{ and } 0.08$) samples prepared by solid-state reaction route followed by spark plasma sintering have been studied. The combination of Pb for Bi substitution and Cu vacancies introduction increases the electrical conductivity while moderately decreasing the Seebeck coefficient. The intrinsically low thermal conductivity of BiCuSeO is noticeably decreased by the introduction of Pb_{Bi} and V_{Cu} of a small concentration ($x = 0.02$) and increased with further increase in x . The optimum balance between the Seebeck coefficient and electrical conductivity was achieved for $x = 0.06$, while for $x > 0.02$, the μ_w/κ_{lat} ratio essentially does not increase and remains at the same level as for $x = 0.02$, limiting the potentially achievable zT_{max} . As a result, the value of zT was eventually increased by more than 3 times compared to the undoped BiCuSeO, and $zT_{max} = 0.75$ at 773 K was achieved for $Bi_{0.94}Pb_{0.06}Cu_{0.94}SeO$. By careful analysis of the literature data and comparing it with our results, we revealed that for $Bi_{1-x}Pb_xCu_{1-x}SeO$ -based oxy-selenides, the main factor determining zT_{max} is the κ_{lat} . We believe that highlighting the mechanisms determining zT value in doped BiCuSeO will help future researchers to form more effective synthesis/doping strategies.

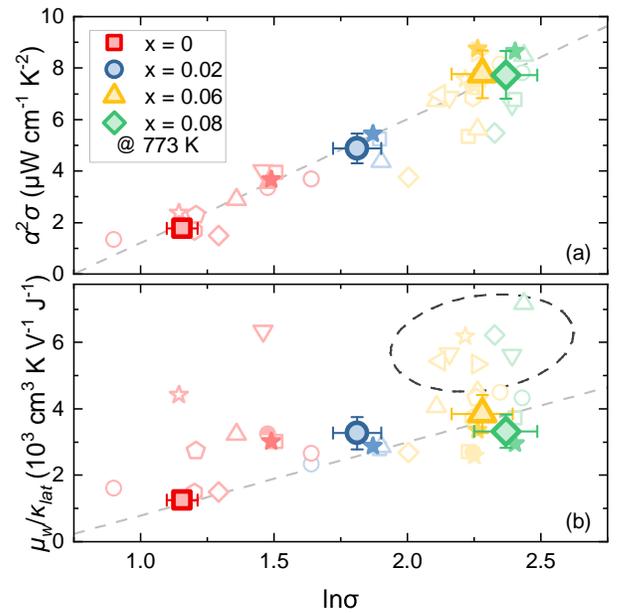


Figure 4 (a) Power factor $\alpha^2\sigma$ and (b) μ_w/κ_{lat} ratio as functions of $\ln\sigma$ for the $Bi_{1-x}Pb_xCu_{1-x}SeO$ ($x = 0, 0.02, 0.06, \text{ and } 0.08$) samples. Literature data for other $Bi_{1-x}Pb_xCu_{1-x}SeO$ -based oxy-selenides are also shown for comparison (Chen et al.,²³ Gu et al.,^{29,30} Lan et al.,¹⁶ Lei et al.,³¹ Li et al.,³² Liang et al.,³³ Liu et al.,²⁶ Pan et al.,³⁴ Ren et al.,^{27,28} Xu et al.,³⁵ Zhu et al.⁹). All displayed data points correspond to values obtained at 773 K.

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CRedit Statement

Aleksandra Khanina: Investigation, Formal analysis, Visualization, Writing – Original draft, Writing – Review & editing; **Andrei Novitskii:** Conceptualization, Formal analysis, Investigation, Methodology, Visualization, Writing – Original draft, Writing – Review & editing; **Daria Pashkova:** Data curation, Investigation, Formal analysis, Visualization; **Andrei Voronin:** Project administration, Funding acquisition; **Takao Mori:** Resources, Funding acquisition, Supervision, Writing – Review & editing. **Vladimir Khovaylo:** Resources, Supervision, Writing – Review & editing.

Conflicts of Interest

There are no conflicts to declare.

Data Availability Statement

Data will be made available on request.

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