

1 **Flaw size-dependent mechanical interlayer coupling and edge-**  
2 **reconstruction embrittlement in van der Waals materials**

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31

32 **Abstract**

33 Van der Waals (vdW) materials with strong in-plane covalent bonding and weak interlayer  
34 interactions consist of two-dimensional (2D) building blocks such as graphene and monolayer  
35 hexagonal boron nitride (*h*-BN). These multilayer 2D materials exhibit exceptional mechanical  
36 characteristics, showcasing impressive fracture resistance as demonstrated in monolayer *h*-  
37 BN. Nonetheless, preserving these remarkable fracture properties in vdW materials remains a  
38 challenge. Here we reveal an anomalous mechanical interlayer coupling that involves  
39 interlayer-friction toughening and edge-reconstruction embrittlement during the fracture of  
40 multilayer *h*-BN. Both asynchronous and synchronous fracture modes and their flaw-size  
41 dependence are identified. Edge reconstruction in the synchronous fracture mode can eliminate  
42 a lattice asymmetry induced toughening mechanism in monolayer *h*-BN leading to  
43 embrittlement of the multilayer *h*-BN, while the asynchronous fracture mode results in greater  
44 fracture resistance. Such findings will provide fundamental guidelines for engineering  
45 interlayer interactions in vdW materials including heterostructures and layered architectures  
46 for better mechanical and functional performances.

47

48 **Main Text**

49 Pristine monolayer two-dimensional (2D) materials have been attracting worldwide  
50 attention over the last two decades owing to their exceptional properties. The relatively weak  
51 van der Waals (vdW) interlayer interactions and the ultrahigh surface-to-volume ratio of  
52 monolayers, allow assembly of monolayers with the same or different chemical compositions  
53 into multilayer structures, opening up the whole new field of vdW materials. Many exotic  
54 behaviors have been demonstrated in these multilayer structures, such as unconventional  
55 superconductivity in bilayer graphene at a magic twisting angle <sup>1</sup>, super-lubricity through  
56 incommensurate contacts between 2D layers <sup>2-4</sup>, formation of moiré excitons that can  
57 resonantly couple with phonons, and interlayer p-n junction as well as interlayer tunneling  
58 effects <sup>5-8</sup>. These interesting interlayer coupling behaviors are thought to be activated by  
59 tailoring the stacking order (AA', AB, etc.), stacking sequence (graphene/h-BN,  
60 MoS<sub>2</sub>/graphene, etc.), and twisting angle, which inspires fundamentally new approaches to  
61 device engineering <sup>1,9-11</sup>. Despite this exciting progress, studies on the effects of interlayer  
62 coupling on the overall mechanical performance, especially fracture behavior, of multilayer  
63 structures are still in a nascent stage.

64 The emerging vdW materials with strong in-plane covalent bonding and weak interlayer  
65 interactions can be constructed from diverse 2D building blocks, and it is well known that these  
66 2D building blocks in their monolayer forms exhibit exceptional mechanical properties such as  
67 ultrahigh strength <sup>12,13</sup>. However, for many practical uses of various vdW materials including  
68 heterostructures and layered architectures, retaining outstanding mechanical properties of  
69 monolayers in their multilayer structures remains a major challenge. One famous example is  
70 the strong and stiff graphene versus the relatively weak and soft graphite. The lack of strong  
71 interlayer interaction is generally believed to be the main contributing factor for such a  
72 challenge. Therefore, understanding, and more importantly, developing corresponding  
73 strategies to maintain high strength and fracture resistance of 2D materials from monolayer to  
74 multilayer vdW materials becomes a very important topic to be explored.

75 It is currently hypothesized that interlayer friction could increase the toughness of  
76 multilayer structures when compared to their monolayer counterparts by retarding both crack  
77 initiation and propagation. This toughening effect has been successfully demonstrated in both  
78 multilayer graphene and MoS<sub>2</sub> systems <sup>14,15</sup>. For example, the strength and elastic modulus of  
79 monolayer graphene is as high as 130 GPa and 1 TPa <sup>12</sup>, respectively. The fracture toughness

80 of monolayer graphene however, is only 4 MPa.m<sup>1/2</sup> <sup>16</sup>, while multilayer graphene's fracture  
81 toughness can reach as high as 8.65±0.46 MPa.m<sup>1/2</sup> due to interlayer-friction toughening <sup>14</sup>.

82 Multilayer *h*-BN (MBN) is an exciting material to study the multilayer effects on fracture  
83 behaviors due to the unique fracture behavior of single layer *h*-BN (SBN). Like graphene, SBN  
84 has a 2D hexagonal lattice, albeit with alternating boron (B) and nitrogen (N) atoms instead of  
85 all carbon (C) atoms, an ultrahigh strength (~100 GPa) and an ultrahigh elastic modulus  
86 (~0.8 TPa) <sup>13</sup>. As the toughest monolayer 2D material verified by both experimental  
87 observations and theoretical predictions, SBN is shown to be intrinsically tough due to its  
88 asymmetric lattice structure. This asymmetric lattice structure exhibits repeated occurrences of  
89 crack deflection and branching along the crack path <sup>13</sup>. Due to the atomic polarity between B  
90 and N atoms, thermodynamically stable MBN generally prefers the AA' stacking order, unlike  
91 multilayer graphene which is in the AB stacking order <sup>17</sup>. MBN has been shown to maintain  
92 the superior strength and elastic modulus from its monolayer constituent <sup>18</sup>, SBN, but there  
93 have been very few quantitative experimental studies on the fracture behaviors of MBN in AA'  
94 stacking order. Earlier fracture studies on MBN with disordered stacking demonstrated that the  
95 fracture properties of MBN do not directly benefit from the effects of interlayer friction or  
96 intralayer lattice asymmetry <sup>19</sup>.

97 In this work, we investigated the fracture behaviors of MBN. We observed two fracture  
98 modes, controlled by the initial flaw size, in AA'-stacked MBN through quantitative *in-situ*  
99 measurements and systematic multiscale mechanics analyses. Depending on the initial flaw  
100 size, AA' MBN exhibited either a toughening effect due to interlayer friction and intralayer  
101 lattice asymmetry, or an embrittlement effect due to interlayer B-N bond formation. This  
102 understanding of MBN fracture could provide important guidance on engineering the interlayer  
103 mechanical coupling in MBN and other vdW materials. In a broader sense, the intriguing  
104 structure-property relationship unveiled in this work could also contribute to new design  
105 principles of layered architectures as already successfully demonstrated in nacre <sup>20,21</sup> and  
106 lamella structures <sup>22,23</sup>.

## 107 **Results**

108 MBN in AA' stacking order was prepared using the mechanical exfoliation method. Details  
109 can be found in **Supplementary Section 1**. A high-quality bulk single crystal *h*-BN was  
110 exfoliated several times and then transferred to a silicon wafer. MBN samples were then  
111 carefully examined under an optical microscope. Samples with the proper size, shape, and

112 thickness, of around 10 nm, were subsequently dry transferred to the Rice push-to-pull (PTP)  
113 mechanical testing devices. The Rice PTP device converts a compressive force on the device  
114 to an in-plane tensile force on a sample. A nanoindenter inside the scanning electron  
115 microscope (SEM) was used to actuate the PTP devices. To measure the sample strain precisely,  
116 Pt fiducial markers were deposited on the sample surface for digital image correlation (DIC)  
117 analyses, as indicated in **Fig. 1**. Three representative *in situ* tensile tests of MBN samples were  
118 selected to demonstrate the intrinsic fracture behaviors of MBN. **Fig. 1 (a)** and **(b)** show the  
119 in-plane tensile test of Sample #1 (MBN ribbon with suspended length  $L = 2.354 \mu\text{m}$ ,  
120 suspended width  $W = 2.804 \mu\text{m}$ , and thickness  $t = 8 \text{ nm}$ ). The tensile strength and failure strain  
121 of Sample #1 is  $\sim 9.48 \text{ GPa}$  and  $\sim 1.81\%$ , respectively. As observed in **Fig.1 (b)**, Sample #1  
122 fractured asymmetrically with cracks deflecting and branching along the rough crack path,  
123 which is commonly seen in its monolayer component, SBN<sup>13</sup>. To measure the fracture  
124 toughness and more closely observe the intrinsic fracture behavior of MBN, a central pre-crack  
125 with a specific length ( $2a$ ) was cut with a focus ion beam (FIB) in Sample #2 and #3, as shown  
126 in **Fig. 1 (c)** and **(e)**.

127 The dimensions of Sample #2 are  $L = 2.6 \mu\text{m}$ ,  $W = 5.36 \mu\text{m}$ , and  $t = 10 \text{ nm}$  with a central  
128 pre-crack ( $2a \sim 0.847 \mu\text{m}$ ), as shown in **Fig. 1 (c)** and **(d)**. The tensile strength and failure strain  
129 of Sample #2 is  $\sim 8.96 \text{ GPa}$  and  $\sim 1.3\%$ , respectively, which are very close to the strength and  
130 failure strain of Sample #1. Interestingly, the crack path leading to final fracture completely  
131 bypassed the central pre-crack and the crack morphology was rough, similar to the crack  
132 morphology observed in MBN without a pre-crack in **Fig. 1 (b)**. Such flaw insensitive (FI)  
133 behavior is quite unexpected in a multilayer 2D material system. According to the Chen and  
134 Gao model for FI behavior<sup>24</sup>, the width of Sample #2 falls into the range of FI due to the  
135 ultrahigh fracture toughness of SBN and possible interlayer-friction toughening in MBN. More  
136 detailed analyses and derivations are presented in the **Supplementary Section 2** and **3**. Owing  
137 to the interlayer-friction toughening effect, the critical energy release rate ( $J_C$ ) of MBN should  
138 be no less than twice of SBN's, i.e.  $172.7 \text{ J/m}^2$ , and the critical SIF ( $K_C$ ) of MBN is no less than  
139  $11.8 \text{ MPa}\cdot\text{m}^{1/2}$  accordingly. The failure of MBN is strength-dominated, and rough fracture  
140 surfaces are present owing to the asymmetrical fracture behaviors.

141 However, as seen in **Fig. 1 (e)** and **(f)**, an interesting fracture mode of Sample #3, which  
142 differed significantly from Sample #2, is unveiled. The size of Sample #3,  $L = 2.742 \mu\text{m}$ ,  $W =$   
143  $5.064 \mu\text{m}$ , and  $t = 9 \text{ nm}$ , is similar to Sample #2, but it has a longer pre-crack length of  $2a =$   
144  $1.65 \mu\text{m}$ . In this case, ideal brittle fracture behavior with a straight crack path and clean

145 cleavage were captured, which is totally different from both SBN<sup>13</sup>, and Samples #1 and #2.  
146 According to the **Supplementary Table S1**, we further categorize the test results into three  
147 groups: MBN with no pre-crack, MBN with a short pre-crack ( $2a < 0.85 \mu\text{m}$ ), and a long pre-  
148 crack ( $2a > 1.2 \mu\text{m}$ ). The summarized failure strengths (solid lines) and predicted limiting  
149 strengths (dashed lines) are shown in **Fig. 1 (g)**. The failure strength of MBN with a short pre-  
150 crack is  $7.62 \pm 0.94 \text{ GPa}$ , which is close to that of MBN with no pre-crack  $7.67 \pm 1.41 \text{ GPa}$ .  
151 Interestingly, the value is even a little higher than its predicted FI strength ( $6.59 \text{ GPa}$ ), which  
152 indicates the limiting strength with FI behaviors as predicted in **Supplementary Section 3**.  
153 However, the failure strength of MBN with a long pre-crack dropped drastically to  $1.91 \pm 0.19$   
154  $\text{GPa}$ , compared to the predicted FI strength  $5.38 \text{ GPa}$ , implying an entirely different failure  
155 mode. Furthermore, the calculated critical stress intensity factor (SIF) of Sample #3  
156 surprisingly dropped to  $K_{\text{IC}} = 1.98 \text{ MPa}\cdot\text{m}^{1/2}$ , which is unexpected based on previous findings  
157 in SBN<sup>13</sup>, and also contradictory to the prediction of Griffith relation  $\sim 3.3 \text{ MPa}\cdot\text{m}^{1/2}$ <sup>13,25</sup>. These  
158 interesting but somewhat contradictory observations imply that flaw size in MBN plays an  
159 important role in triggering significant changes in the fracture behavior of this multilayer  
160 system.

161 To reveal the underlying role played by flaw size in the fracture of MBN, three typical  
162 samples with pre-crack lengths of  $2a = \sim 0.8 \mu\text{m}$ ,  $1.2 \mu\text{m}$  and  $1.4 \mu\text{m}$  are shown in **Fig. 2**. As  
163 shown in **Fig. 2 (a)**, the failure strength and strain is reduced with increasing pre-crack length  
164 as expected, but the measured critical SIF ( $K_{\text{IC}}$ ) counterintuitively shows the same trend which  
165 normally should be independent of pre-crack length<sup>13,26,27</sup>. To better observe the fracture edge  
166 morphology and understand the underlying fracture mechanisms, TEM characterization was  
167 carefully performed on fractured MBN samples. PTP devices with fractured samples were  
168 directly loaded into the TEM to avoid any potential damage to the fracture edges. Crack  
169 deflections were clearly observed in samples with short pre-crack length, as shown in **Fig. 2**  
170 **(b)** and **(c)**, which is consistent with the asymmetric fracture characteristics of SBN<sup>13</sup>. The  
171 characteristics of asynchronous fracture featuring stair-case edges are shown in **Fig. 2 (c)**,  
172 which is commonly seen in AB stacked multilayer 2D systems<sup>14,15</sup>. The rough fracture edges  
173 were clearly observed at both the micro- and nanometer length scales. To highlight  
174 asynchronously fractured layers with rough fracture edges, the TEM image is colored-coded  
175 based on contrast difference as shown in **Fig. 2 (c)**. As the flaw size (pre-crack length) in MBN  
176 increases, the asynchronous fracture behavior fades away at the micrometer scale and no stair-  
177 case crack edges are observed from SEM images as shown in **Fig. 2 (d)**. However, from the

178 TEM image shown in **Fig. 2 (e)**, rough fracture edges and stair-case crack edges still exist at  
179 the nanometer scale, indicating that asynchronous fracture persists to a limited extent.  
180 Therefore, the asynchronous fracture behavior, represented in the form of the average spacing  
181 between crack edges of different layers, diminishes from several hundreds of nanometers to a  
182 few nanometers, as shown in **Fig. 2 (d)** and **(e)**. The asynchronous fracture mode is further  
183 suppressed with the increase of the pre-crack length. As shown in **Fig. 2 (f)** and **(g)**, atomically  
184 smooth fracture edges were observed at both the micro- and nanometer length scales. A  
185 distinguishing feature observed is a straight dark line, which is consistent with the  
186 characterization of closed edges of multilayer 2D materials<sup>28,29</sup>. The insert in **Fig. 2 (g)** shows  
187 the closed edge structures. According to the TEM simulations, it is reasonable to have a darker  
188 contrast for the closed edges due to the higher local atomic density. Details can be found in  
189 **Supplementary Section 1**. It can be inferred that only simultaneous fracture along the same  
190 crack path in different layers could result in such uniform and smooth fracture edges as shown  
191 in **Fig. 2 (g)**, demonstrating a very intriguing synchronous fracture process.

192 To confirm this highly interesting flaw-size dependent fracture behavior, a series of tests  
193 on MBN samples with different pre-crack lengths were conducted. A summary of the measured  
194 or estimated (in flaw-insensitive cases) fracture toughness data as a function of the pre-crack  
195 length was listed in **Supplementary Table S1**. Samples with FI fracture behaviors were  
196 commonly found when pre-crack length was below 0.85  $\mu\text{m}$ . With the increase of the pre-crack  
197 length, the FI fracture behaviors became absent, and the average fracture toughness quickly  
198 decreases to  $\sim 2 \text{ MPa}\cdot\text{m}^{1/2}$ . Meanwhile, the fracture modes of MBN were observed to transition  
199 from asynchronous to synchronous fracture.

200 Based on *in situ* quantitative fracture tests and high-quality SEM and TEM images, it is  
201 clearly seen that there exist two very different fracture modes in MBN: asynchronous fracture  
202 vs. synchronous fracture. Asynchronous fracture mode (fracture occurs along different crack  
203 paths in each layer) generally leads to a higher fracture resistance and rough (staircase) fracture  
204 edges, while synchronous fracture mode (fracture appears to occur along a common crack path  
205 in each layer) leads to a reduced fracture resistance and ultra-smooth fracture edges. It appears  
206 that the initial flaw (pre-crack) size in MBN plays an important role in determining the  
207 transition between these two fracture modes (**Fig. 2**), but the physical origin and underlying  
208 mechanisms remain elusive. Therefore, multiscale mechanics analyses were performed to  
209 understand the structural origins of dual fracture modes in MBN and to explore the underlying  
210 mechanisms associated with the observed flaw-size dependent fracture behaviors.

211 The unique AA' stacking order of MBN makes it, under certain circumstances, favor  
212 synchronous fracture because of the possible formation of interlayer B-N bonds on fracture  
213 edges of adjacent layers leading to straight and closed edges. To get a complete picture of such  
214 behavior, we analyzed two distinctive fracture processes in MBN: (I) intralayer B-N bond  
215 breakage without interlayer B-N bond formation; (II) intralayer B-N bond breakage with  
216 interlayer B-N bond formation. In fracture process (II), the interlayer B-N bond formation leads  
217 to edge reconstruction. Fully reconstructed edges can have two possible patterns of bi-planar  
218 edge reconstruction — type I edge reconstruction and type II edge reconstruction — both of  
219 which are studied carefully (**Fig. 3**).

220 A multiscale mechanics analysis under ideal  $K_I$  loading field is carried out to investigate  
221 the intrinsic fracture behaviors of MBN during the crack growth initiation stage as shown in  
222 **Fig. 3 (a)**. The yellow region is modelled by the *ab initio* method to accurately capture the  
223 atomic reconstruction in the fracture processing zone. The  $J$  integral in **Fig. 3 (b-e)** is only an  
224 evaluation of local loading to record the corresponding critical events at crack tip, but not a  
225 prediction for energy release rate, considering the lattice trapping effect<sup>30</sup>. Detailed  
226 descriptions of the analysis are presented in the **Supplementary Section 4**. The accurate  $J_C$   
227 and  $K_C$  of MBN under the fracture process (I) and (II) follows the analyses in the  
228 **Supplementary Section 2** and **5**. Following the fracture process (I), if none of the fractured  
229 edges behind the crack tip are reconstructed, the interlayer interactions of MBN are dominated  
230 by interlayer friction, which shows strong asynchronous characteristics due to the lack of out-  
231 of-plane mirror symmetry in MBN. As shown in **Fig. 3 (b)**, the crack tip bifurcates when  $J =$   
232  $13.0 \text{ J/m}^2$ , close to that of SBN, which bifurcates at  $J = 11.5 \text{ J/m}^2$ <sup>13</sup>. Then the bifurcated tip  
233 fails at  $J_C = 17.5 \text{ J/m}^2$ , featuring an asynchronous characteristic — cracks on adjacent layers  
234 deflect in different directions. Following the fracture process (II), if interlayer B-N bonds at  
235 fractured edges behind the crack trip are formed and lead to edge reconstruction, the formation  
236 of interlayer bond will dominate the interlayer interaction of MBN, which results in  
237 synchronous fracture because edge reconstruction stitches the two exposed fractured edges  
238 together into a folded structure thereby reducing the fracture energy. According to the Griffith  
239 relation, the accurate  $J_C$  and  $K_C$  of MBN with complete edge reconstruction can be obtained by  
240 evaluating the formation energy of MBN's folding structure in unit length,  $J_C = \sim 3.80 \text{ J/m}^2$  and  
241  $K_C = \sim 1.75 \text{ MPa}\cdot\text{m}^{1/2}$  (See **Supplementary Section 5**). The fracture of MBN in process (II) is  
242 synchronous across adjacent layers, leading to brittle fracture with a lower fracture toughness  
243 as observed experimentally. More details about crack growth in adjacent layers can be found

244 in **Fig. 3 (d)** and **(e)**. **Fig. 3 (d)** presents typical asynchronous fracture behaviors. Once the  
245 crack initiates, crack bifurcation and repeated deflection occur in different directions of  
246 adjacent layers due to the lack of both in-plane and out-of-plane mirror symmetry in MBN as  
247 shown in **Fig. 3 (d)**. MBN with asynchronous fracture demonstrates higher fracture toughness  
248 than SBN due to the interlayer-friction shielding effect <sup>13</sup>. **Fig. 3 (e)** presents typical  
249 synchronous fracture behaviors. The crack grows straight in-plane with edge reconstructions  
250 following the crack tip. Interestingly, benefitting from the unique AA' stacking order, there  
251 can be two possible types of edge reconstruction in MBN as shown in **Fig. 3 (f)** and **(g)**. For  
252 type I edge reconstruction, interlayer B-N bond formation occurs between the same two layers  
253 on both side of the crack; while for type II edge reconstruction, individual layers form interlayer  
254 B-N bond with a different layer on each side of the crack. In both cases, the crack tip advances  
255 in a typical brittle manner and leaves a straight crack path. In the synchronous fracture mode,  
256 no dangling bond is formed from crack propagation, and the MBN exhibits a much lower  
257 fracture toughness than SBN <sup>13</sup>. In other words, the in-plane lattice asymmetric toughening  
258 mechanism that made SBN the toughest 2D material <sup>13</sup> can be diminished in MBN due to edge  
259 reconstructions following the crack tip.

260 Based on our experimental observations of the flaw-size dependent fracture behaviors in  
261 MBN, it is reasonable to assume that the formation of out-of-plane B-N bonds in edge  
262 reconstructions is also flaw-size dependent. Here, two models are formulated to comprehend  
263 this flaw-size dependence employing multiscale mechanics analyses as shown in **Fig. 4 (a)**. In  
264 **Fig. 4 (a)**, the MBN sample under uniaxial tension along the  $y$  direction is prepared with a  
265 central pre-crack along the  $x$  direction. The in-plane boundary of the continuum region is  
266 defined by the analytical solution of the displacement field containing a central pre-crack under  
267 uniaxial loading <sup>31</sup>. The out-of-plane boundary is periodic. Detailed settings of the model are  
268 presented in the **Supplementary Section 6**.

269 The first model is to unveil the relationship between the mobility of the reconstructed  
270 crack tip and the flaw size in MBN, as illustrated in **Fig. 4 (b)**. When a reconstructed crack tip  
271 advances lattice by lattice, no dangling bond could be formed, and the translational symmetry  
272 is preserved. Therefore, based on this feature, a stress-augmented activation model is proposed  
273 to clarify the relationship between the mobility of the reconstructed crack tip and the flaw size  
274 in MBN under critical loading <sup>32,33</sup>. The flaw size is obtained from the measured pre-crack  
275 length of MBN samples tested via *in situ* tensile experiments. The MBN samples with varying  
276 flaw sizes are then categorized into two groups: failure with FI behaviors (denoted as FI and

277 colored red) and failure without FI behaviors (denoted as non-FI and colored blue). The flaw-  
278 size dependence of fracture behaviors can be extracted from the predicted mobility of the  
279 reconstructed crack tip. Detailed derivations can be found in the **Supplementary Section 7**.  
280 As shown in **Fig. 4 (c)**, it is obvious that the propagation of the reconstructed crack tip will be  
281 more favorable as the flaw size grows longer. When the flaw size is smaller than 0.8  $\mu\text{m}$ , the  
282 mobility of reconstructed crack tip is reduced to less than  $1 \times 10^{-10}$  nm/s, which is unlikely to be  
283 detected under normal conditions, implying that the reconstructed crack tip can hardly  
284 propagate owing to the suppressed edge reconstruction. As the flaw size grows larger than 1.0  
285  $\mu\text{m}$ , the mobility of reconstructed crack tip is significantly increased according to the stress-  
286 augmented activation model, eventually giving rise to the synchronous fracture mode.

287 The second model is to explore the relation between the energy landscape in edge  
288 reconstruction and the flaw size in MBN. As shown in **Fig. 4 (b)**, the edge reconstruction in  
289 MBN can be decomposed into two events, intralayer cleavage and interlayer healing. The  
290 intralayer cleavage is denoted as the breakage of the  $\text{B}^1\text{-N}^1$  bond and the  $\text{N}^2\text{-B}^2$  bond, and the  
291 interlayer healing is the formation of the  $\text{B}^1\text{-N}^2$  bond and the  $\text{B}^2\text{-N}^1$  bond. The energy consumed  
292 by intralayer cleavage is a constant, theoretically twice of the B-N bond formation energy in a  
293 bilayer *h*-BN model, none of which is dependent on flaw-size. Meanwhile the energy barrier  
294 of interlayer healing, the final event in edge reconstruction shown in **Fig. 4 (d)**, varies due to  
295 the intralayer and interlayer interatomic distance, which is affected by the in-plane flaw-size  
296 dependent boundary conditions. Therefore, to explore the flaw-size dependent energy  
297 landscape of edge reconstruction is to quantify the relationship between the energy barrier of  
298 the interlayer healing and the flaw size. The energy barrier can serve as a measure of the  
299 capability for the edge reconstruction to follow the crack tip closely during crack propagation,  
300 which is a suitable criterion to predict whether the crack propagation is synchronous or  
301 asynchronous. Obviously, an easy-to-overcome energy barrier will allow edge reconstruction  
302 to keep up with the crack tip lattice by lattice and lead to a synchronous fracture under critical  
303 loading. The basic setting for calculating the reconstructed energy barrier is based on the MBN  
304 model with a central pre-crack in **Fig.4 (a)**. By changing the central pre-crack length and  
305 extracting the reconstructed energy barrier, the flaw-size dependence of the reconstructed  
306 energy barrier could be assessed. Usually, a volume change, due to atomic reconstruction while  
307 investigating the energy barrier of a phase transition in a finite cell composed of several atoms,  
308 could complicate the results severely. Therefore, the generalized solid-state nudged elastic  
309 band (G-SSNEB) method, which is widely adopted in exploring the energy barrier of phase

310 transitions in a small system<sup>34</sup>, is utilized in the *ab initio* modeling to comprehensively evaluate  
311 the atomic reconstruction aligned with the corresponding cell volume change in the *z* dimension  
312 (out-of-plane direction with periodic boundary condition). Details can be found in  
313 **Supplementary Section 8**. Clear flaw-size dependence of the energy barrier  $\Delta E_b$  in interlayer  
314 healing is shown in **Fig. 4 (e)**, no matter which type of edge reconstruction occurs as indicated  
315 in **Fig. 4 (d)**. The  $\Delta E_b$  decreases with increasing flaw size, indicating that edge reconstruction  
316 occurs much more easily in the presence of a longer pre-crack, which confirms the unique flaw-  
317 size dependent dual fracture behaviors observed in MBN. Due to the experimental challenges  
318 of directly observing the interlayer bond formation in our samples, the proposed interlayer  
319 bonding in the synchronized fracture mode remains to be verified by experiments.

## 320 **Conclusion**

321 It is of significant importance to understand how to preserve outstanding properties of  
322 monolayer constituents in multilayer assemblies. The fascinating fracture behaviors of SBN  
323 and the unique structure of MBN provide an ideal test bed for exploring the influence of  
324 interlayer mechanical coupling on the fracture behaviors of multilayer nanomaterials. In MBN,  
325 interlayer friction associated with the vdW interaction can be treated as a type of weak  
326 interlayer mechanical coupling, which can lead to flaw-insensitive and asynchronous fracture  
327 behaviors. On the other hand, activated by out-of-plane B-N bond formation at the fractured  
328 edges in adjacent layers, strong interlayer mechanical coupling enabled by edge reconstructions  
329 could result in flaw-sensitive and synchronous fracture behaviors. By increasing the pre-crack  
330 (flaw) size in MBN, the transition from asynchronous fracture to synchronous fracture was  
331 observed experimentally and explained by multiscale mechanics analyses. With the rise of vdW  
332 solids or 2.5D materials<sup>35</sup>, solid materials and structures constructed by engineering interlayer  
333 interactions represent an important paradigm shift in designing novel materials for many  
334 exciting applications<sup>4,36-38</sup>. The findings from this study represent a solid step forward towards  
335 designing and engineering more robust vdW materials against fracture and inspiring the  
336 optimal design of layered architectures.

337

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343 affiliation of Prof. Huajian Gao has changed to Mechano-X Institute, Applied Mechanics  
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346

### 347 **Author Contributions**

348 Z. S., B. Z. and Y. Y. conceived the project. B. Z., Y. Y. and J. L. designed the experiments  
349 and analyzed the data. Z. S. and H. G. performed simulations and analyzed the results. G. G.,  
350 Y. H. and D. T. performed TEM analysis. Q. F., B. Shin., D. S., Q. A. and X. Z. helped with  
351 the sample preparation and characterization. Y. Z. performed digital image correlation. T. T.  
352 and K. W. performed material synthesis. Z. S. B. Z. Y.Y. supervised by N. P., B. Sheldon., J.  
353 L. and H. G. drafted the manuscript with inputs, discussion, and approval from all co-authors.  
354 Z. S., B. Z. and Y.Y. contributed equally to this work.

355

### 356 **Competing Interests**

357 The authors declare no competing interests.

358

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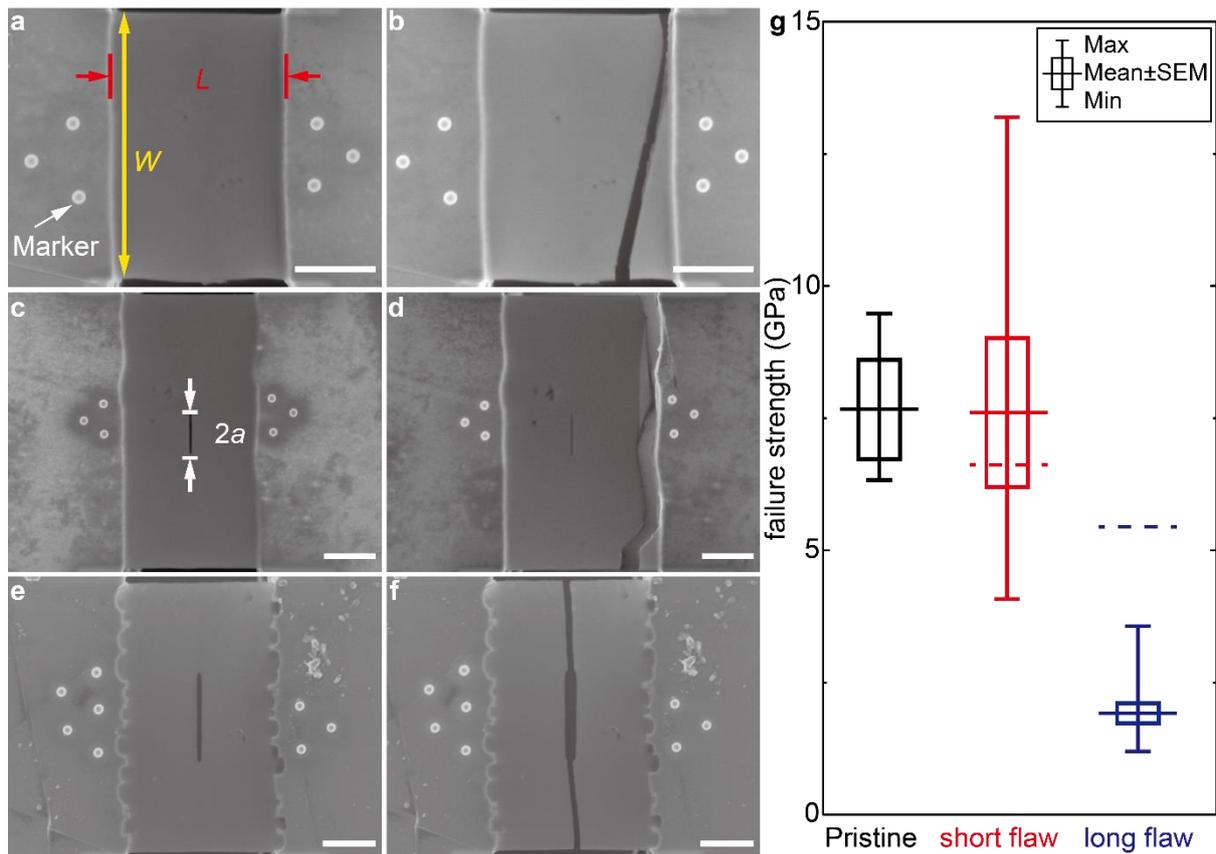
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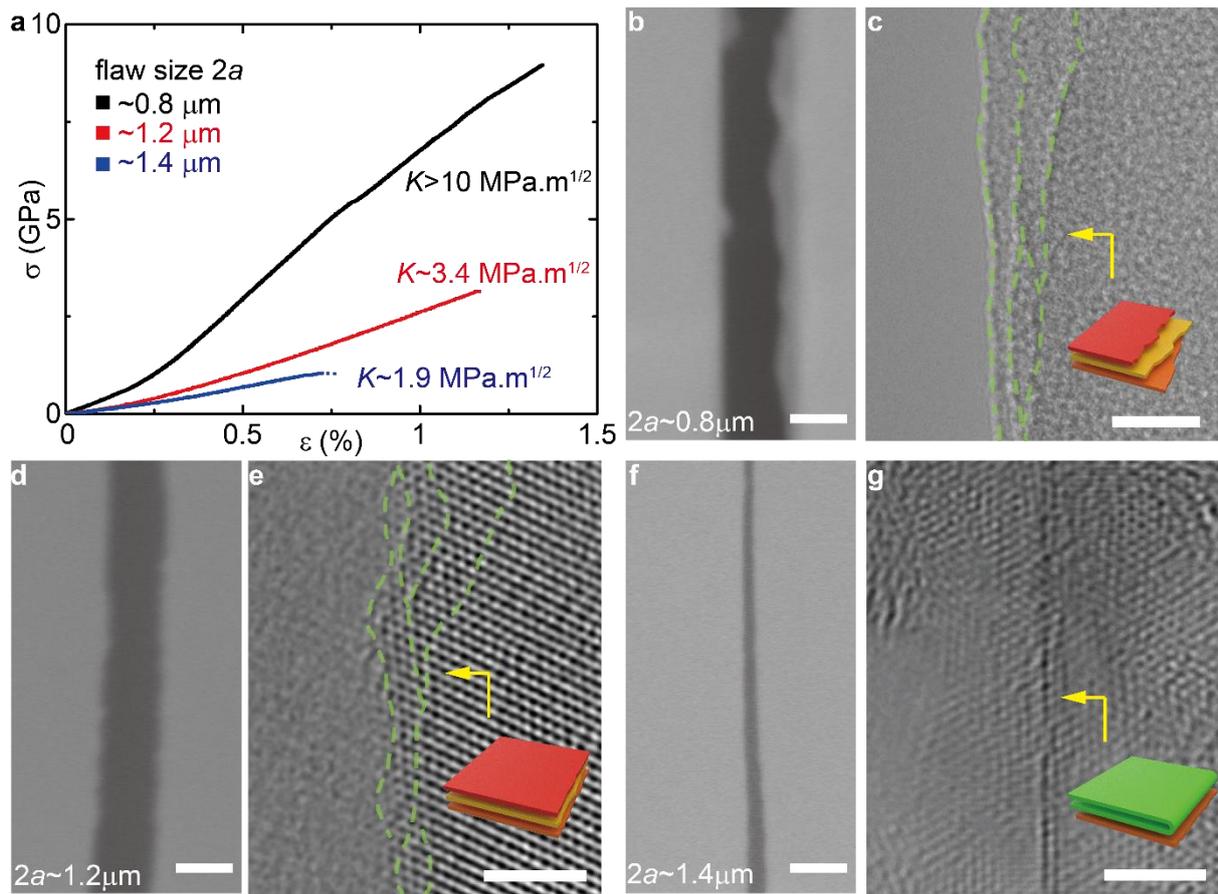
448

449 **Fig. 1. In-situ quantitative tensile tests of MBN.** (a) Suspended MBN without pre-crack  
 450 before a tensile test, and (b) after a tensile test; (c) Flaw insensitivity observed in a suspended  
 451 MBN with a short pre-crack before a tensile test, and (d) after a tensile test; (e) Flaw sensitivity  
 452 observed in a suspended MBN with a long pre-crack before a tensile test, and (f) after a tensile  
 453 test; (g) The failure strength of measured MBN strip. Data are presented as mean values  $\pm$   
 454 SEM, and error bars indicate the range (minima to maxima) of the data ( $n_{\text{pristine}} = 3$ ,  $n_{\text{short flaw}} =$   
 455  $6$ , and  $n_{\text{long flaw}} = 13$ ). Dashed line is the mean of predicted FI strength. Scale bars:  $1 \mu\text{m}$  for (a),  
 456 (b), (c), (d), (e), and (f).

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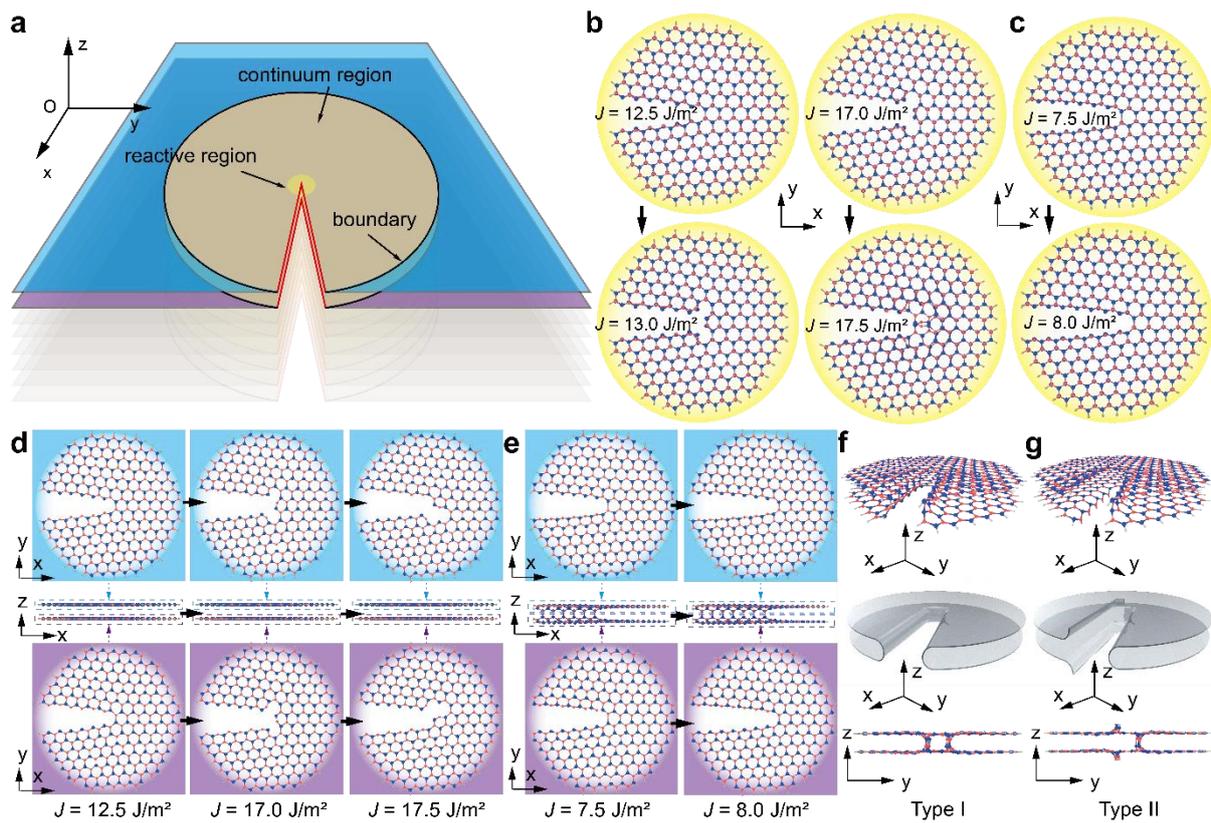
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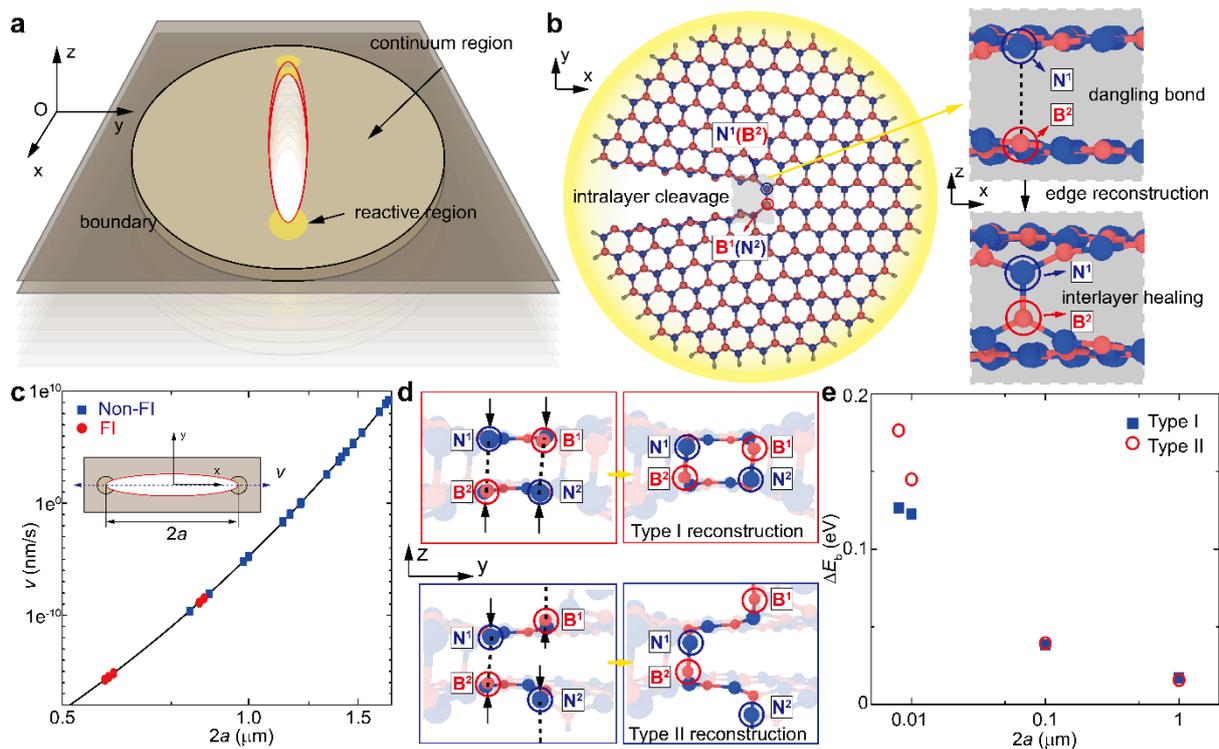
461 **Fig. 2. MBN crack edge characterization after tensile testing.** (a) Stress-strain curves of  
 462 three typical MBN with flaw size  $2a = \sim 0.8 \mu\text{m}$ ,  $\sim 1.2 \mu\text{m}$  and  $\sim 1.4 \mu\text{m}$ ; (b) SEM image of an  
 463 asynchronous crack edge of a fractured MBN with pre-crack length of  $\sim 0.8 \mu\text{m}$  and estimated  
 464 fracture toughness above  $10 \text{ MPa}\cdot\text{m}^{1/2}$ ; (c) TEM images of the rough crack edge, where the  
 465 insert in (c) shows a schematic of the asynchronous fracture edge; (d) SEM image of a straight  
 466 crack edge of a fractured MBN with pre-crack length of  $\sim 1.2 \mu\text{m}$  and fracture toughness around  
 467  $3.4 \text{ MPa}\cdot\text{m}^{1/2}$ ; (e) TEM images of the rough crack edge at atomic scale, where the insert in (e)  
 468 shows a schematic of the transition from asynchronous to synchronous fracture edges; (f) SEM  
 469 image of a fractured MBN with pre-crack length of  $\sim 1.4 \mu\text{m}$  and fracture toughness around  $1.9$   
 470  $\text{MPa}\cdot\text{m}^{1/2}$ ; (g) TEM images of atomically smooth crack edge, where the insert in (g) shows a  
 471 schematic of the synchronous fracture edge. Scale bars: 100 nm for (b), (d), and (f); 20 nm for  
 472 (c); 2 nm for (e) and (g).

473



474

475 **Fig. 3. Ideal  $K_I$  field analysis of dual fracture modes in MBN.** (a) Multiscale schematic of  
 476 crack initiation in MBN; (b) Sequential critical events of crack growth initiation in non-  
 477 reconstructed MBN; (c) Critical events of crack growth initiation in reconstructed MBN; (d)  
 478 Sequential critical events of crack growth initiation in adjacent layers of non-reconstructed  
 479 MBN; The top and bottom layers shown are highlighted by the blue and purple backgrounds.  
 480 (e) Critical events of crack growth initiation in adjacent layers of reconstructed MBN; The top  
 481 and bottom layers shown are highlighted by the blue and purple backgrounds. (f) Type I edge  
 482 reconstruction in MBN; (g) Type II edge reconstruction in MBN.



483

484 **Fig. 4. Flaw-size dependence of fracture modes in MBN.** (a) Multiscale schematic  
 485 illustration of crack propagation in MBN with a central crack; (b) Sequential critical events of  
 486 crack growth initiation in reconstructed MBN (from intralayer cleavage to interlayer healing);  
 487 (c) Theoretical prediction of the relationship between the mobility  $v$  of reconstructed crack tip  
 488 and the flaw size  $2a$ ; (d) Interlayer healing transition in type I and type II reconstructed MBN;  
 489 (e) Relationship between the reconstructed energy barrier  $\Delta E_b$  and the flaw size  $2a$ .

490

491

492 **Methods**

493 **MBN testing sample preparation.** MBN nanosheet was mechanical exfoliated from a bulk *h*-  
494 BN crystals synthesized by the high-pressure and high-temperature method using wafer tape  
495 (Nitto Denko), and was deposited on a clean SiO<sub>2</sub>/Si substrate (UniversityWafer). PMMA  
496 (MicroChem, A4) was spin-coated on the sample substrate and then lifted off from the substrate  
497 using NaOH solution (Sigma-Aldrich). The PMMA/*h*-BN thin film was further transfer to the  
498 copper glider Aperture Grids (Ted Pella). A probe station and a micromanipulator  
499 (Micromanipulator) were used to dry transfer the target MBN sample to the Rice push-to-pull  
500 (PTP) device (**Supplementary Section 1**). Then the PMMA layer was removed through  
501 annealing method in a vacuum tube furnace (MTI corporation, OTF-1200X). After the  
502 annealing process, the samples were carefully screened using SEM for subsequent tests.

503

504 **In-situ SEM mechanical testing.** *In-situ* mechanical test was carried out in a SEM/FIB (FEI,  
505 Helios Nanolab 660) equipped with an in-SEM nanoindentation system (Bruker PI-85L). The  
506 PTP device with the MBN sample loaded on it was fixed on the sample holder using the SEM  
507 sample mounting hard wax (Ted Pella) and a small amount of silver paste (Electron  
508 Microscopy Sciences). Before the mechanical test, FIB was used to define the testing geometry  
509 and manufacture the pre-crack. Then, the nanoindenter tip would push the top beam of the  
510 device, which converts the compression force to the uniaxial tensile force applied to the  
511 samples. The nano-indenter tip was controlled to move at a speed of 10 nm·s<sup>-1</sup>. The quantitative  
512 nanoindentation system recorded the force and displacement of the tip and the deformation and  
513 fracture process were recorded by the SEM system for further analysis and digital image  
514 correlation (**Supplementary Section 1**). Atomic force microscopy (Park, NX20) was used  
515 precisely measure the thickness of the tested MBN sample. The methods used to calculate  
516 measured stress  $\sigma_m$ , strain  $\varepsilon_m$ , stress intensity factor  $K$  and typical results can be found in  
517 **Supplementary Section 1, Fig. S2 and Table S1.**

518

519 **TEM tests and simulation.** After mechanical test, fractured MBN samples on PTP were  
520 directly loaded in the TEM (FEI, Titan). High-resolution transmission electron microscopy  
521 (HRTEM) images are performed with a double-corrected Titan cubed Themis G2 operated at  
522 300 kV. The microscope is equipped with a Ceta camera, a Gatan Quantum 966 energy filter  
523 and an electron monochromator. QSTEM code<sup>39</sup> was used to simulate the TEM images of  
524 bilayer *h*-BN with closed and open edges. Under an accelerating voltage of 300 kV, a spherical  
525 aberration coefficient ( $C_s$ ) of 0.001 mm, and Scherzer defocus condition, the double layer *h*-

526 BN sheets could be treated as weak-phase objects, and intensity of TEM images are  
527 proportional to the atomic potential and phase of electron exit waves. The simulated HRTEM  
528 images were further used to compare with the experimental results. Details can be found in  
529 **Supplementary Section 1.**

530

531 **DFT calculation and vdW correction.** For an accurate description of structural evolution in  
532 *h*-BN, spin-polarized electronic structure optimizations based on the density functional theory  
533 (DFT) are carried out by Vienna Ab-initio Simulation Package (VASP)<sup>40,41</sup>. The generalized  
534 gradient approximation (GGA) in the Perdew Burke Ernzerhof (PBE) version is utilized for  
535 the exchange-correlation functional<sup>42,43</sup>. To avoid the artificial interactions between the 2D  
536 sheet and its periodic images, vacuum space larger than 10 Å is included in the non-periodic  
537 direction. Several vdW schemes have been chosen to calibrate the interlayer interaction of  
538 MBN as shown in **Table S2**. The vdW-D4 is selected for all MBN calculations as the predicted  
539 interlayer distance is the closest to experiment<sup>44,45</sup>.

540

541 **Evaluating the reconstructed energy barrier.** The generalized solid-state nudged elastic  
542 band (G-SSNEB) is used to evaluate the energy barrier of interlayer healing aligned with  
543 corresponding cell volume change in *z* dimension (out-of-plane direction with periodic  
544 boundary condition), which is widely adopted in exploring the energy barrier of phase  
545 transitions in small system<sup>34</sup>. In the *ab initio* G-SSNEB calculations, the initial state is chosen  
546 as the optimized crack tip configuration before intralayer cleavage, while the final state is  
547 defined as the optimized crack tip configuration after interlayer healing. Both of them have  
548 been optimized by conjugate gradient (CG) method. The stress tensor section in the VTST  
549 code<sup>46,47</sup> has been modified in aim to only allow the volume change in *z* direction during the  
550 exploration of energy landscape.

551

## 552 **Data Availability**

553 The data and code that support the findings of this study are available from the corresponding  
554 authors on reasonable request.

555

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557

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