

Supplemental Material for “Velocity-Independent Dry Friction on Mica: Realization of Ideal Amontons-Coulomb Friction”

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This file includes supplemental materials for the detail in experimental conditions, relationship between shear stress and normal stress, friction coefficients of mica at 65 and 100°C, the results of slide-hold-slide tests at 22 and 200°C, available temperature dependence of the “*a*” parameter for various materials, and the detail in molecular dynamics simulations.

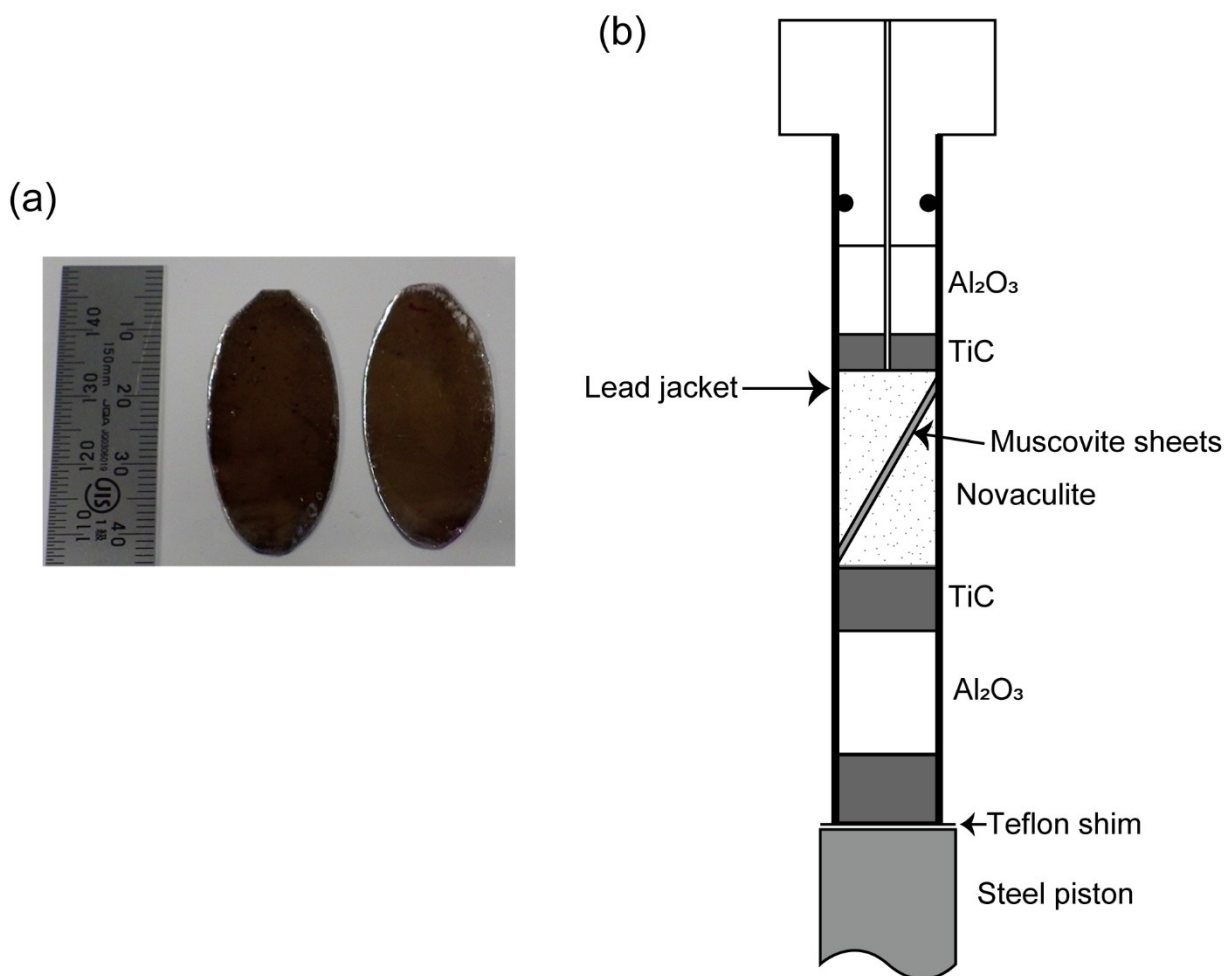


FIG S1. (a) Single-crystal mica sheets. (b) Cross section of the sample assembly with mica sheets between 30° sawcut novaculite driving blocks. During the experiments, the sample column was continuously evacuated through a 1.6 mm-diameter axial hole that extends to the upper porous novaculite driving block.

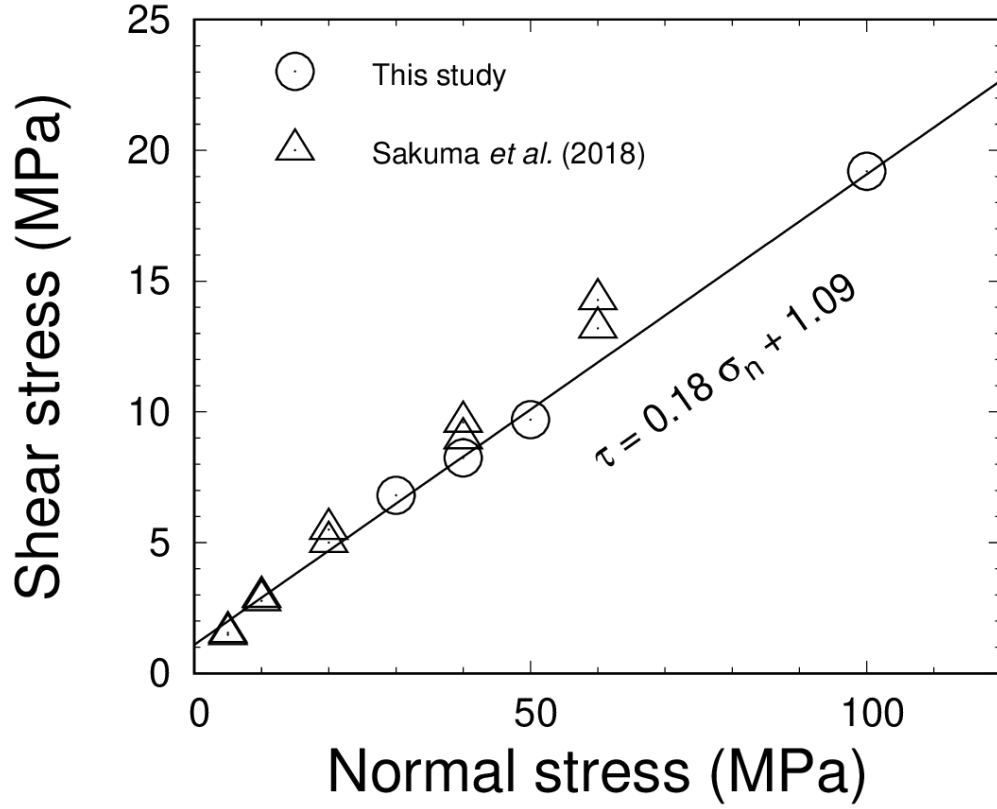


Fig. S2 Relationship between shear stress and normal stress of single-crystal mica measured using a triaxial shear apparatus in this study and a double-direct shear apparatus in a previous study [21] under room-temperature and dry conditions. The equation shown in the figure represents a linear fit to the data obtained in this study.

TABLE S1. Conditions of triaxial shear experiments and obtained “*a*” parameter.

Sample no.	Temperature (°C)	Normal stress (MPa)	Experimental type	<i>a</i>
HS11	22	100	Velocity-stepping	0.0058 (0.0011)
HS14	65	100	Velocity-stepping	0.0028 (0.0005)
HS13	100	100	Velocity-stepping	0.0005 (0.0002)
HS12	200	100	Velocity-stepping	<0.0001 (0.0002)
HS15	22	100	Slide-hold-slide	-
HS16	200	100	Slide-hold-slide	-

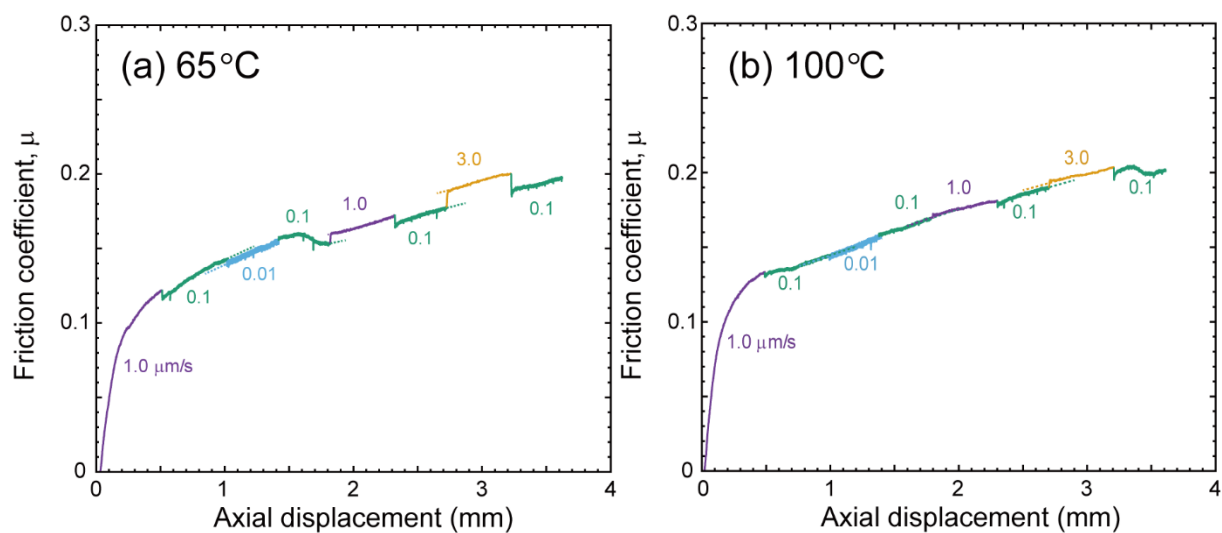


FIG S3. Friction coefficients of mica obtained from velocity-step tests at 65°C (a) and 100°C (b). The numbers near the lines indicate the sliding velocity parallel to the shear planes ($\mu\text{m/s}$). Dotted lines indicate the baseline of the steady-state friction used to calculate the difference in the friction coefficients before and after the velocity change.

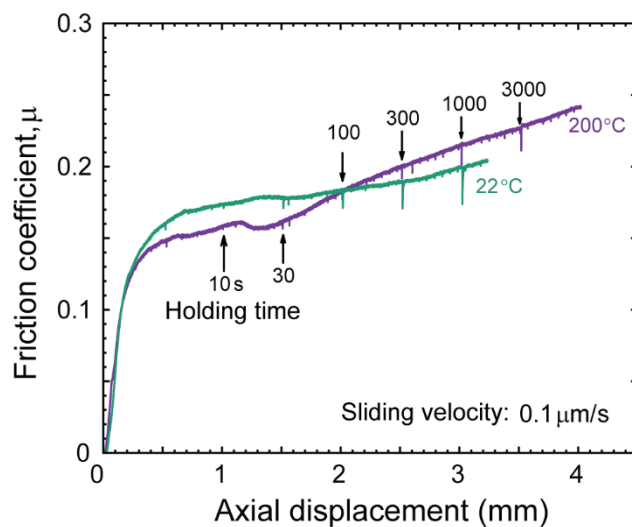


FIG S4. Friction coefficients of mica obtained from slide-hold-slide tests at 22°C and 200°C.

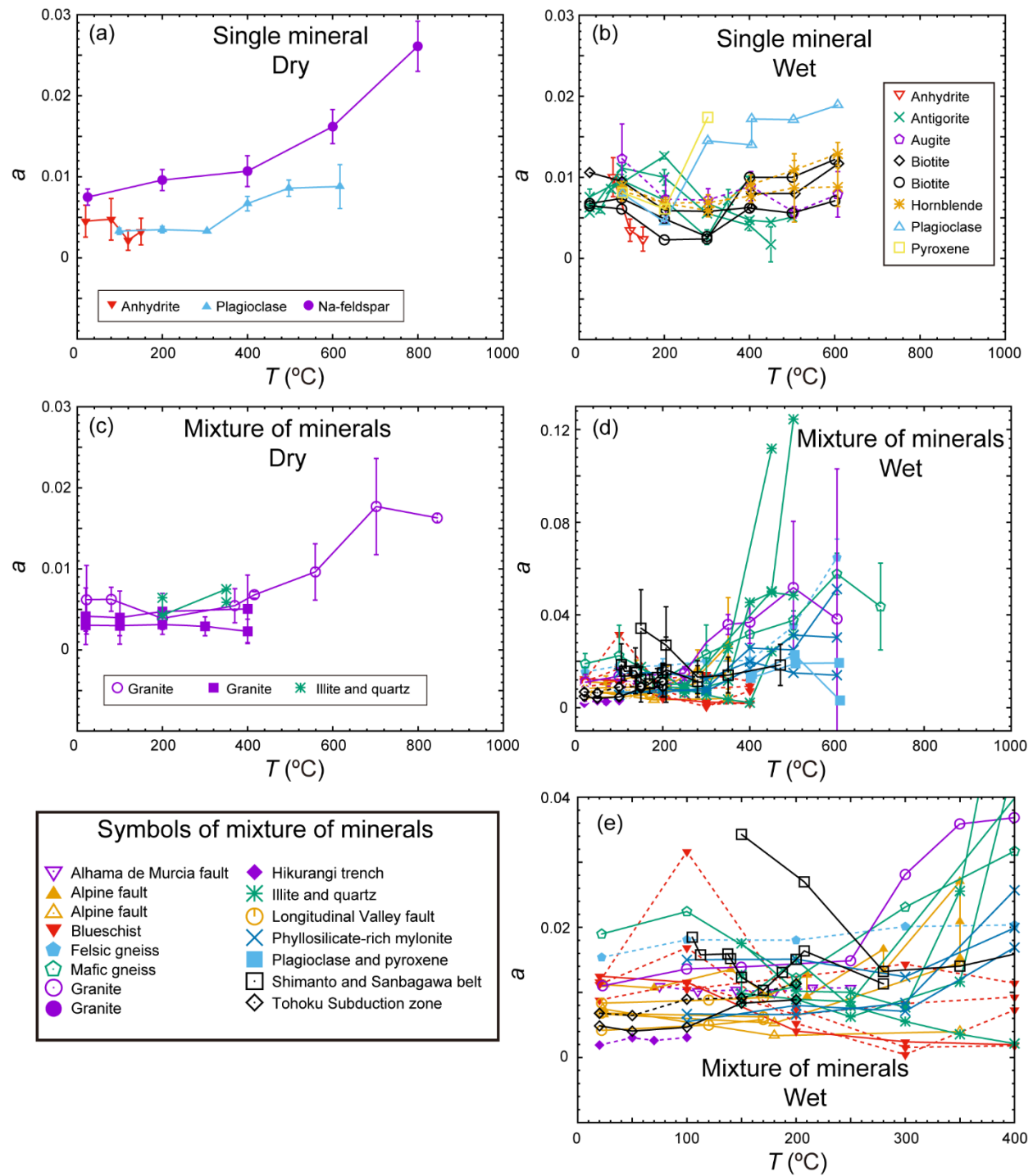


FIG S5. Available temperature dependence of the “ a ” parameter for simulated gouges composed of a single mineral under dry (a) and wet (b) conditions, and a mixture of minerals under dry (c) and wet (d and e) conditions. The data was cited from the literature (anhydrite [35], plagioclase [36,37], Na-feldspar [31], antigorite [38], augite [39], biotite [40,41], hornblende [42], granite [16,43], illite and quartz [44], Alhama de Murcia fault [45], Alpine fault [46,47], blueschist [48], felsic gneiss [49], mafic gneiss [49], Hikurangi trench [50], Longitudinal Valley fault [51], phyllosilicate-rich mylonite [52], Shimanto and Sanbagawa belt [53], Tohoku subduction zone [54]).

Non equilibrium molecular dynamics (NEMD) simulations—Classical molecular dynamics simulations were conducted by using in-house program, MXDTRICL. Muscovite is a 2:1 type sheet-silicate mineral, consisting of an $[\text{AlO}_6]$ dioctahedral sheet sandwiched between two $[(\text{Si},\text{Al})\text{O}_4]$ tetrahedral sheets. One-fourth of the Si^{4+} ions are replaced by Al^{3+} ions in the tetrahedral sheets and the negative charge of the sheets is compensated by the interlayer K^+ ions. Force field of muscovite used in this study is an interatomic potential model, and it has been developed to reproduce the muscovite structure and physical properties [32-34]. The calculated lattice parameters of $2M_1$ muscovite by this model are $a = 5.159 \text{ \AA}$, $b = 8.966 \text{ \AA}$, $c = 20.065 \text{ \AA}$, and $\beta = 95.3^\circ$ at 27°C and 0.1 MPa . The electrostatic Coulomb interaction was calculated by using the Ewald method [55]. The differential equation of motion was solved by the velocity-Verlet algorithm with a time increment of $0.4 \text{ femtosecond (fs)}$. Equilibrium lattice constants were obtained by *NPT* ensemble simulations at pressures of 6 GPa and temperatures of 27 , 100 , and 200°C . The equilibrium lattice constants were used for initial structures for subsequent shear simulations. Dislocation-free and ripplocation-bearing supercells at 27°C before the shear were $a = 25.48 \text{ \AA}$, $b = 26.48 \text{ \AA}$, $c = 19.19 \text{ \AA}$, $\beta = 95.1^\circ$, and $a = 167.1 \text{ \AA}$, $b = 26.35 \text{ \AA}$, $c = 39.58 \text{ \AA}$, $\beta = 94.9^\circ$, respectively. Two ripplocations were included in the ripplocation-bearing supercell; they were built with the length of 34 units of a in the primitive cell on two straight layers with the length of $165 \text{ \AA}$ in a direction (32 units of a in the primitive cell). By applying hydrostatic pressure, the initial structure for shear simulations was obtained. Shear direction and plane were fixed to $\langle 100 \rangle$ and (001) , respectively. Shear strain was applied by tilting the c -axis of the supercell with the shear rates of 10^9 , 10^{10} , and 10^{11} s^{-1} . Friction coefficients as a function of strain were calculated by simply dividing the shear stress by normal stress.